



Jet Modification as a function of energy lost and jet mass depletion

Abhijit Majumder
Wayne State University

Riken BNL High p_T workshop April 2016

Outline

- 1) MC generation and base mode
- 2) Review of MATTER and role of virtuality
- 3) Mass, virtuality scale
- 4) virtuality vs. distance
- 5) From virtuality to reconstructed mass
- 6) Sensitivity to R
- 7) Mass as a window to strong and weak coupling

other messages in this talk!

- Realism!
Calculations on a hydro
- Event generators
the method of the future
- Mixture of strong and weak coupling
a requirement of all calculations.

Motivation

We now have intra-jet data

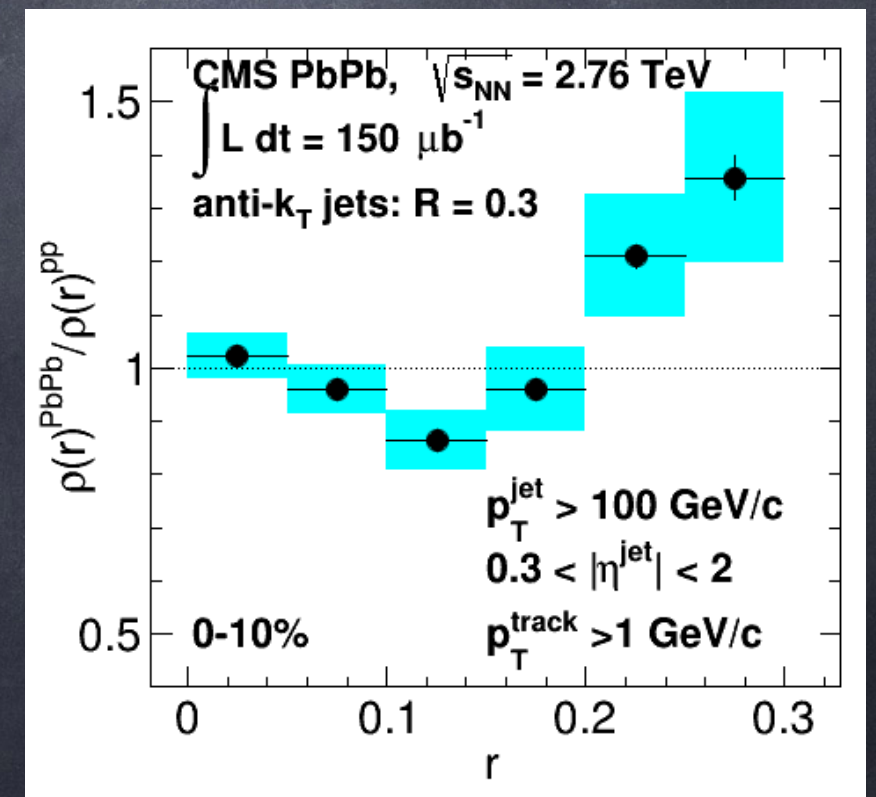
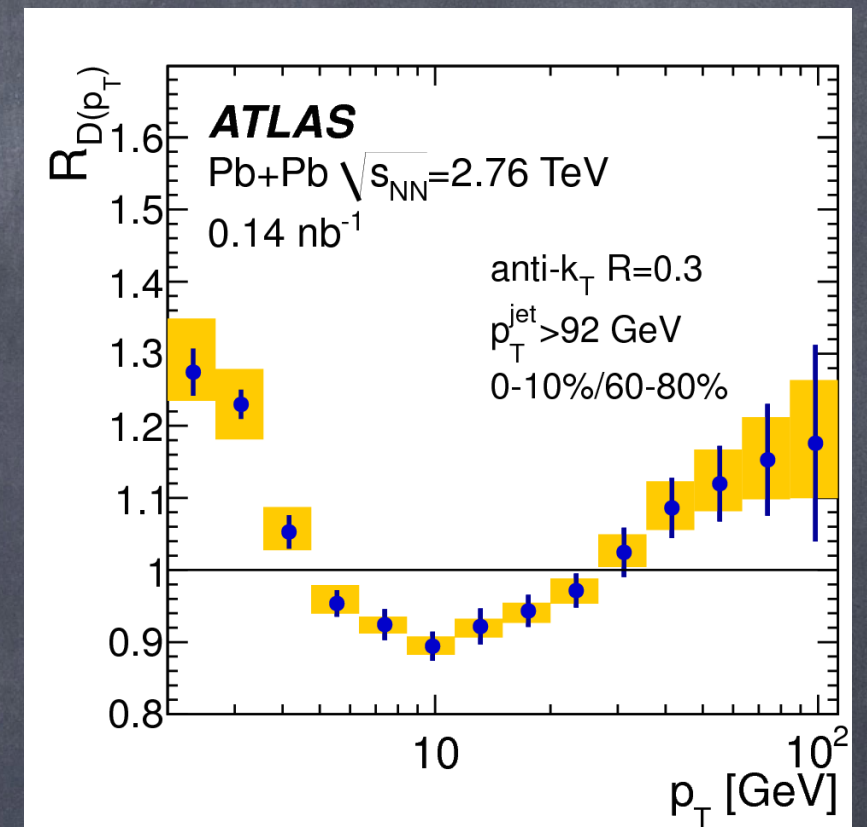
obtained by comparing medium modified and vacuum jets

at the same energy

But is Energy the only way to classify a jet?

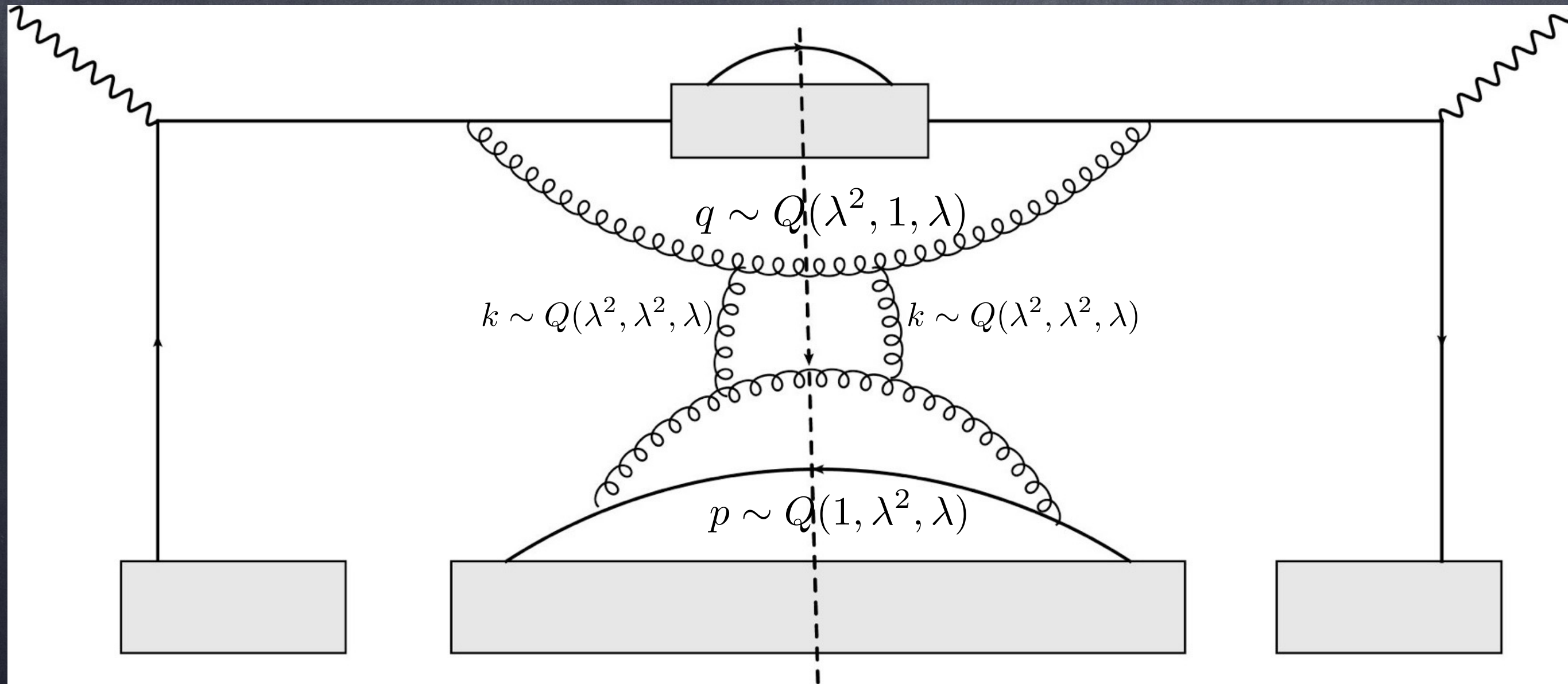
Jet Mass: a boost invariant property of a jet.

To study this, need a MC generator



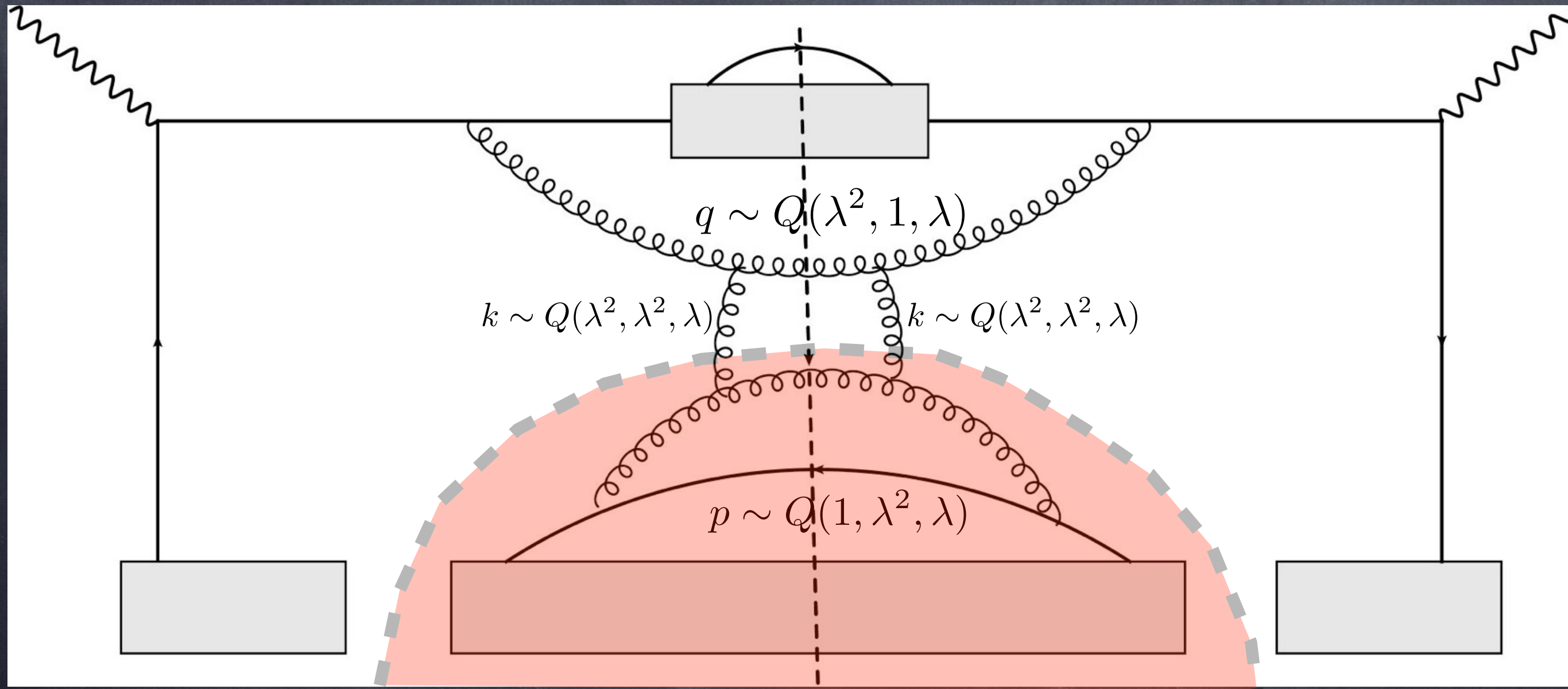
Details of HT-MC code MATTER

Modular All Twist Transverse Elastic scattering and Radiation



Details of HT-MC code MATTER

Modular All Twist Transverse Elastic scattering and Radiation



All of this evolution is hiding within \hat{q}

Light quark modification is sensitive to the high Q^2 , low- x part of the in-medium gluon distribution.

Varieties of MC generators

Bottom-up MC generators:

JEWEL, YaJEM

Top-Down generators:

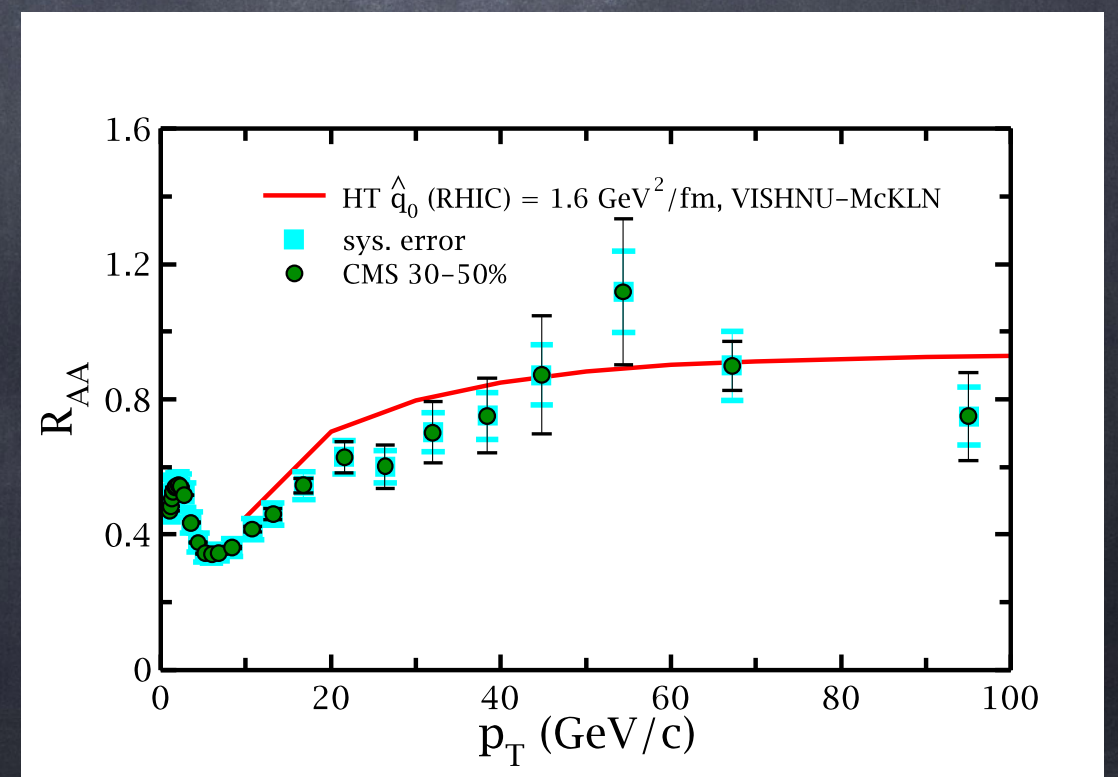
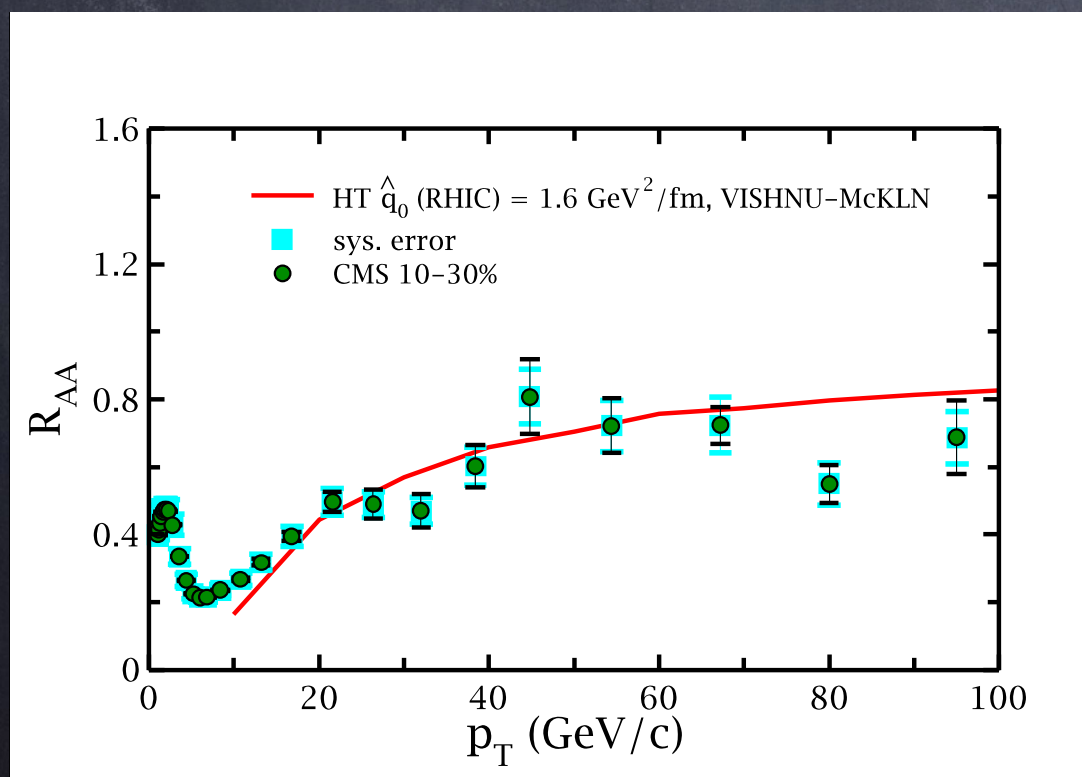
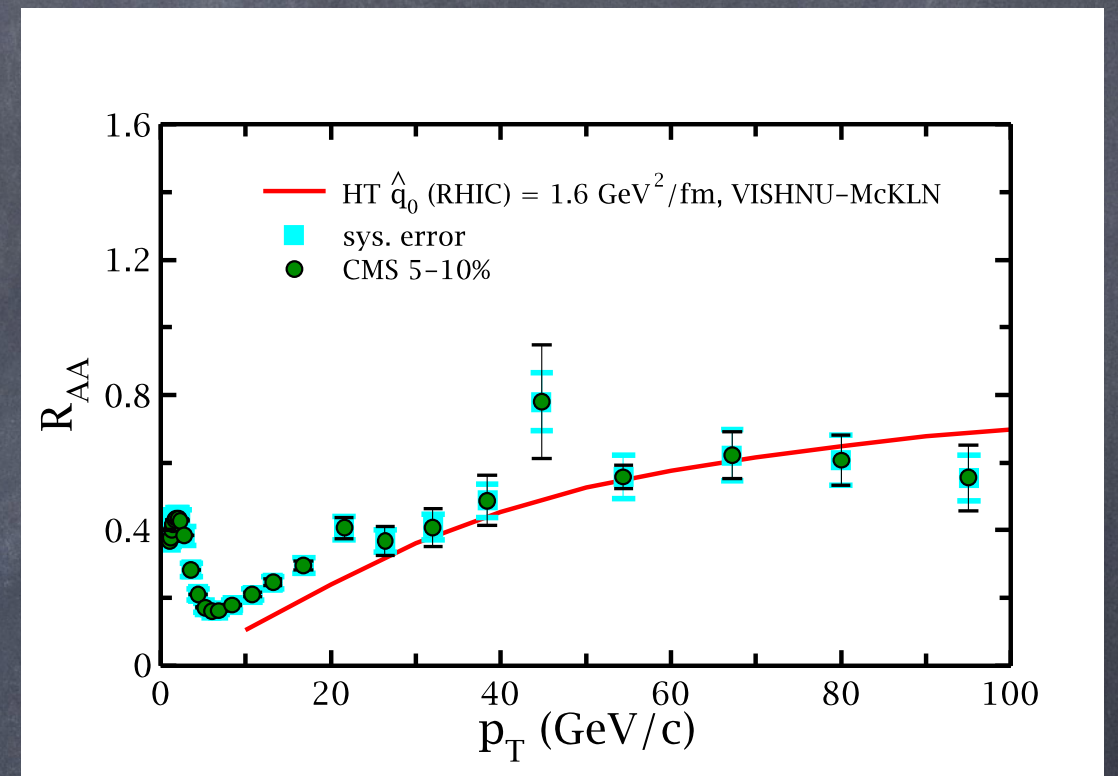
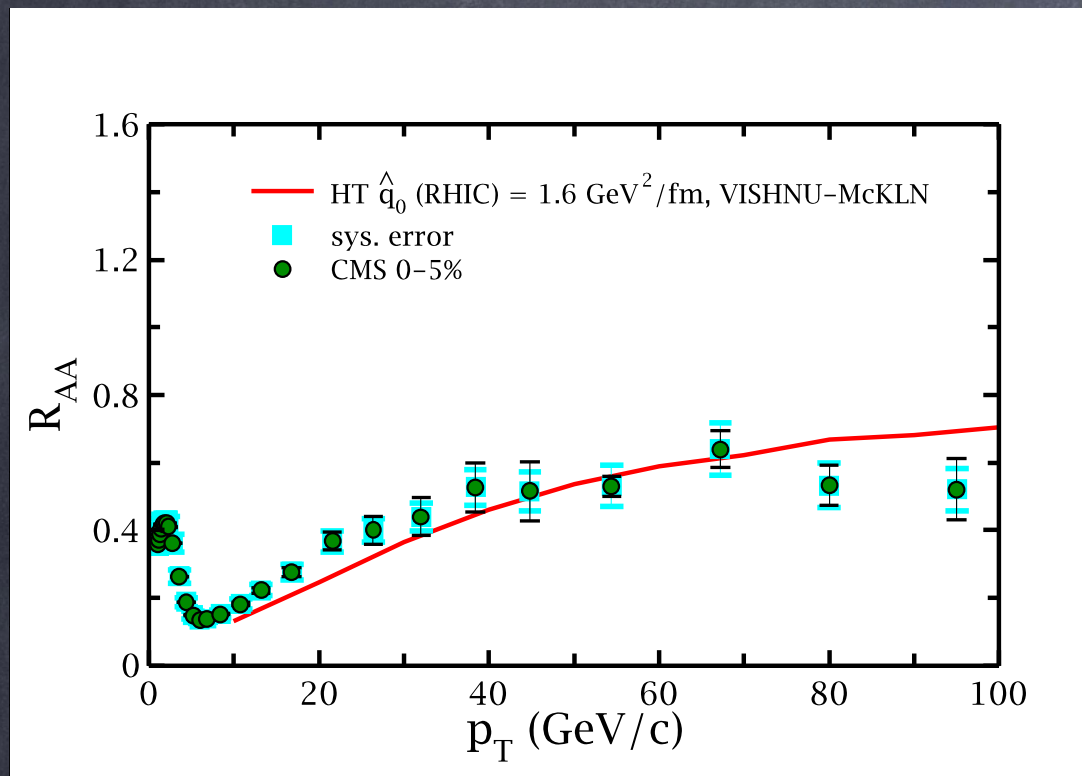
MARTINI, MATTER, Q-PYTHIA

Hybrid generators,

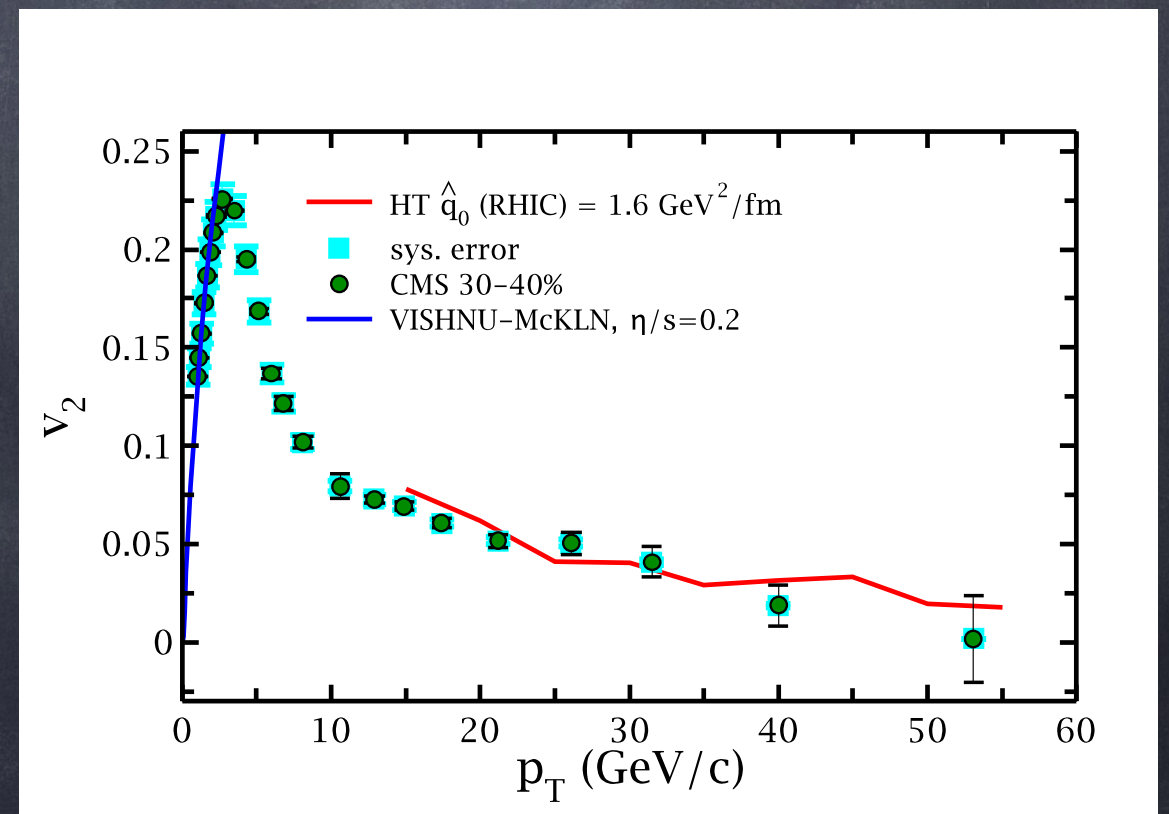
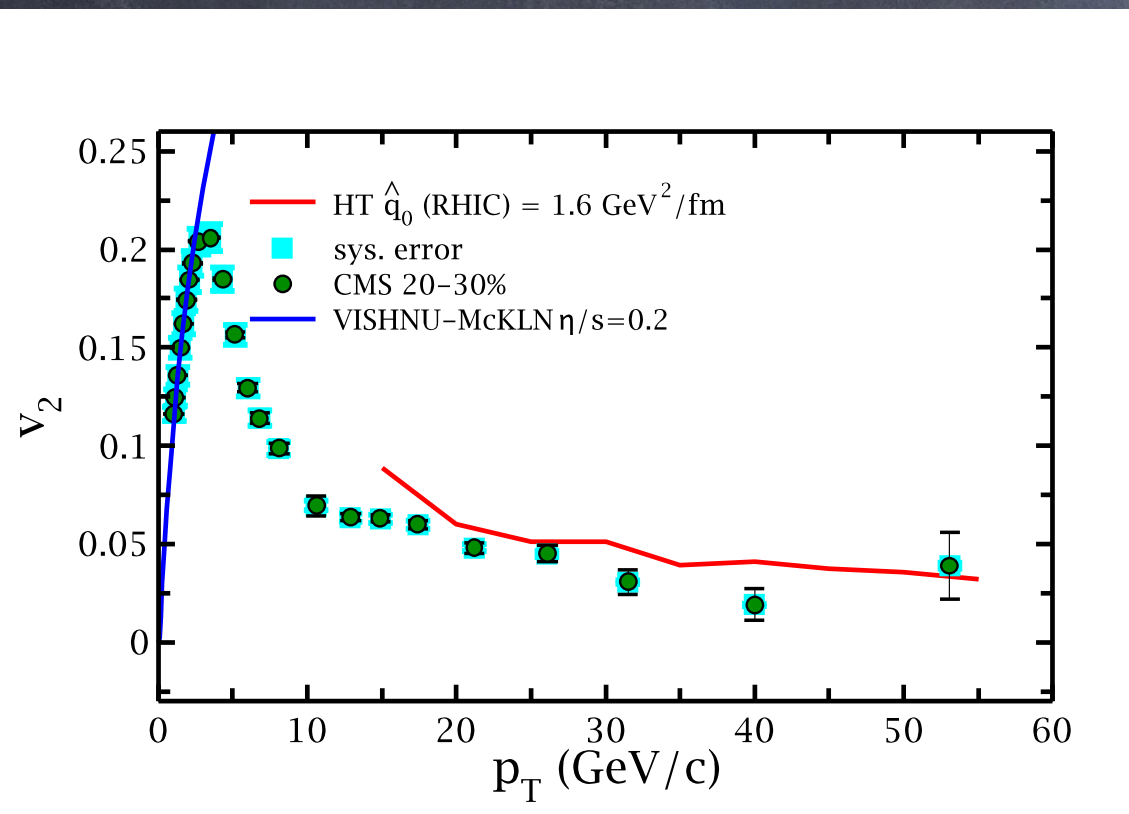
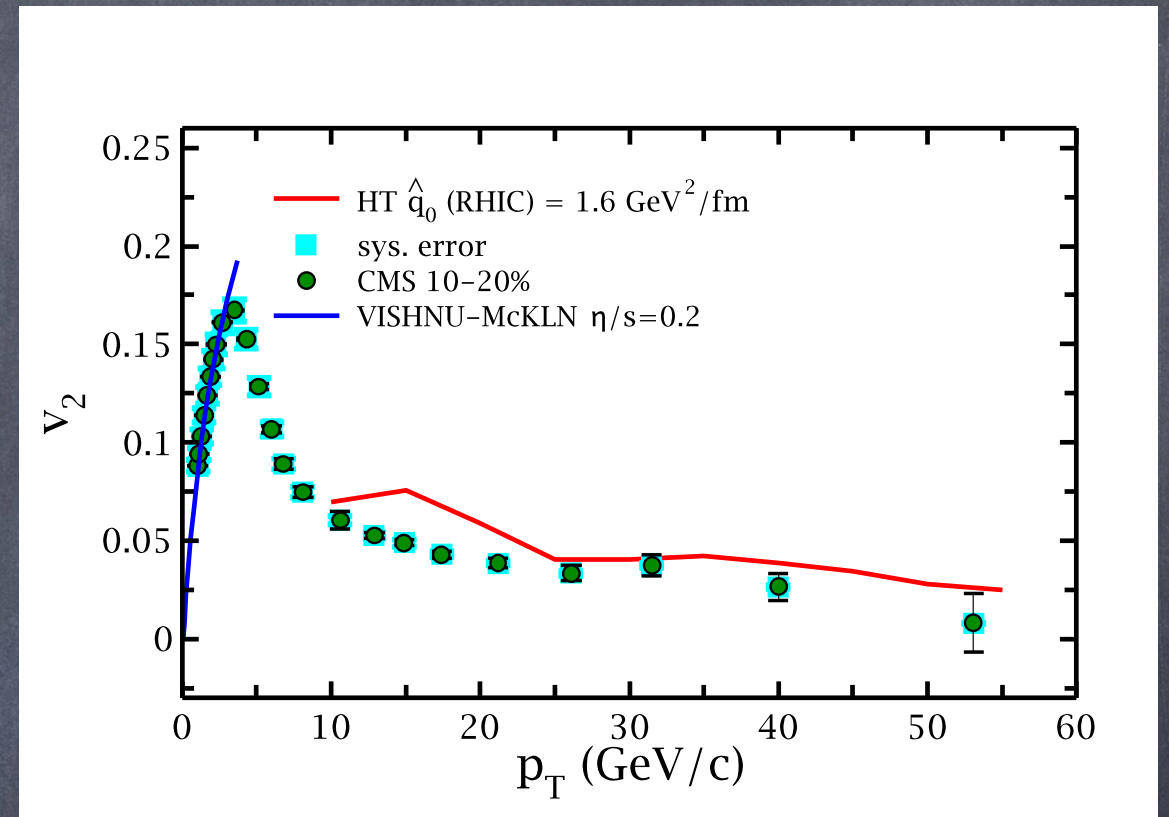
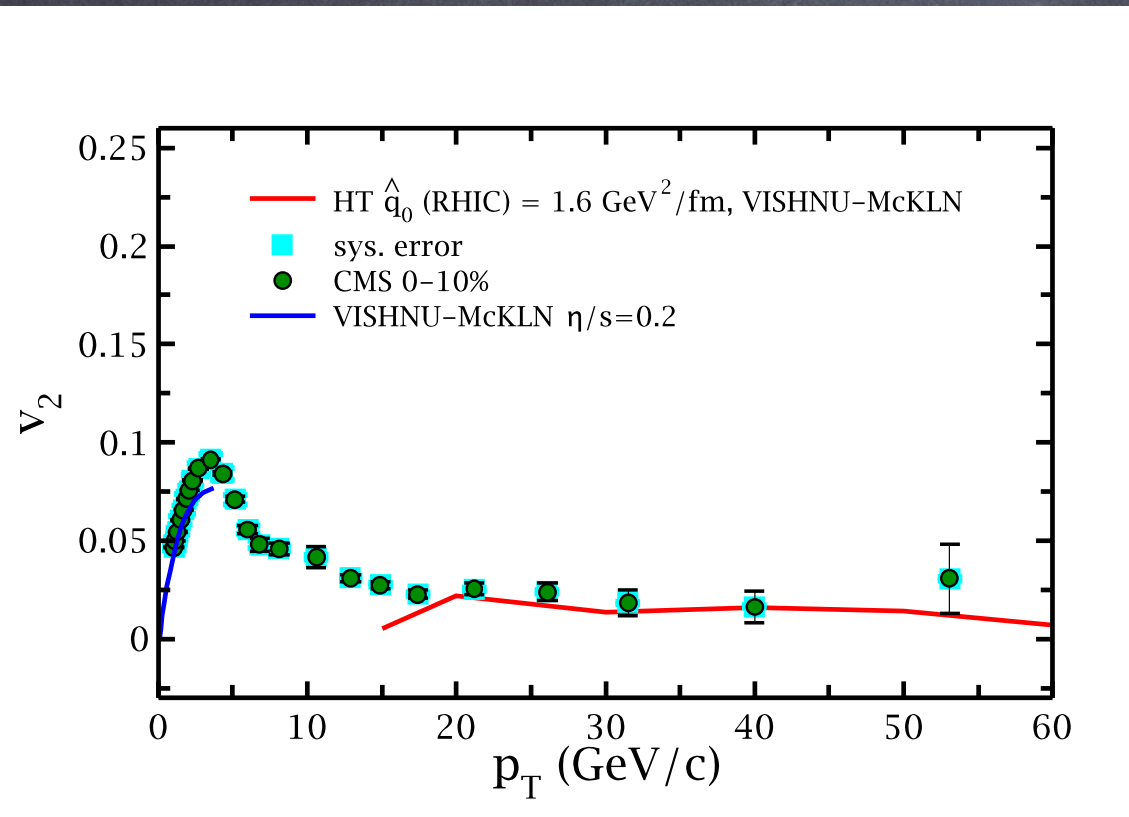
PYQUEN

How good is your base formalism?

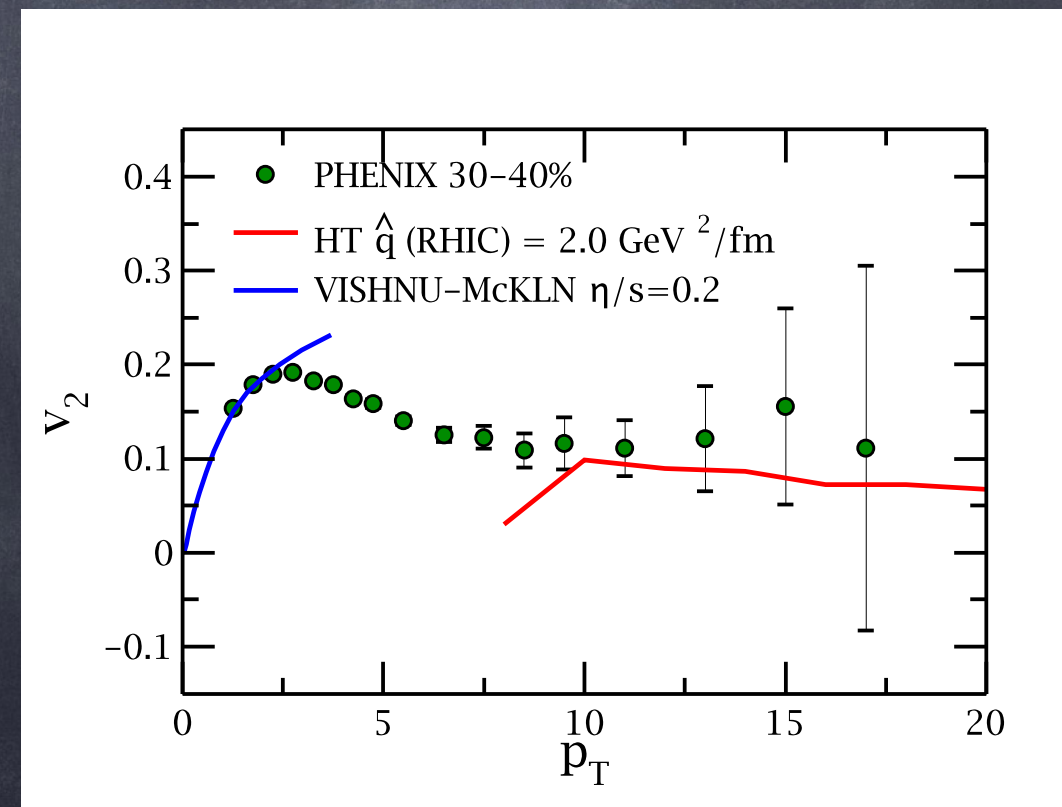
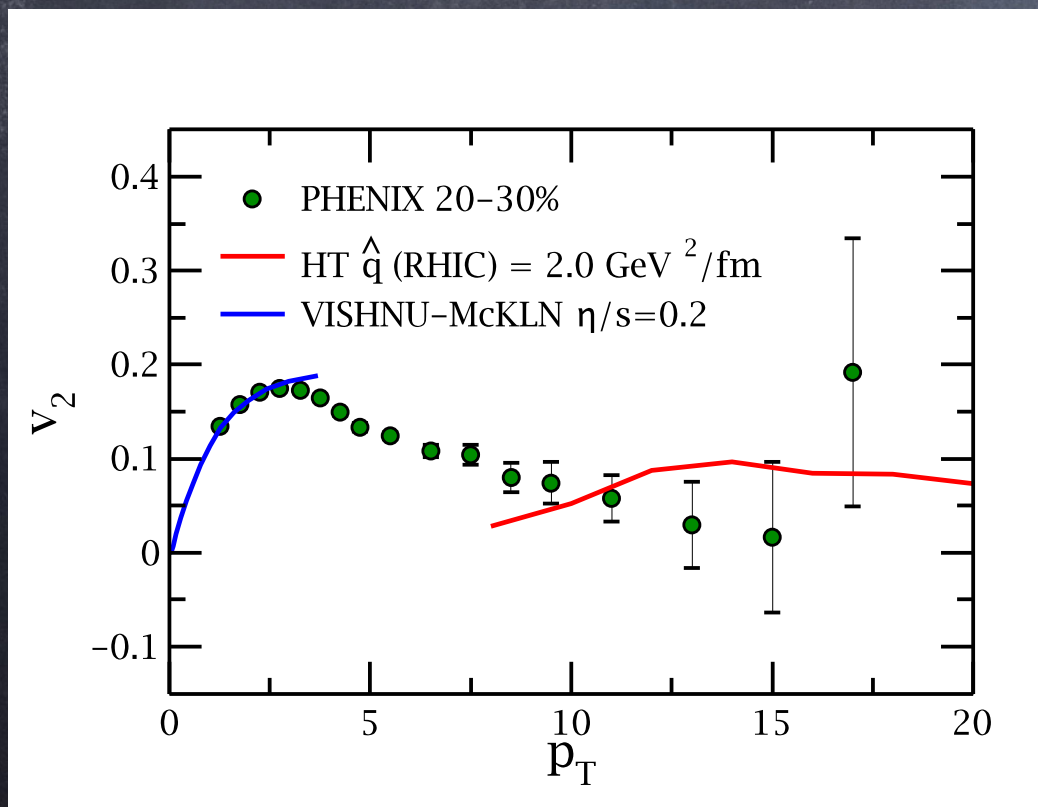
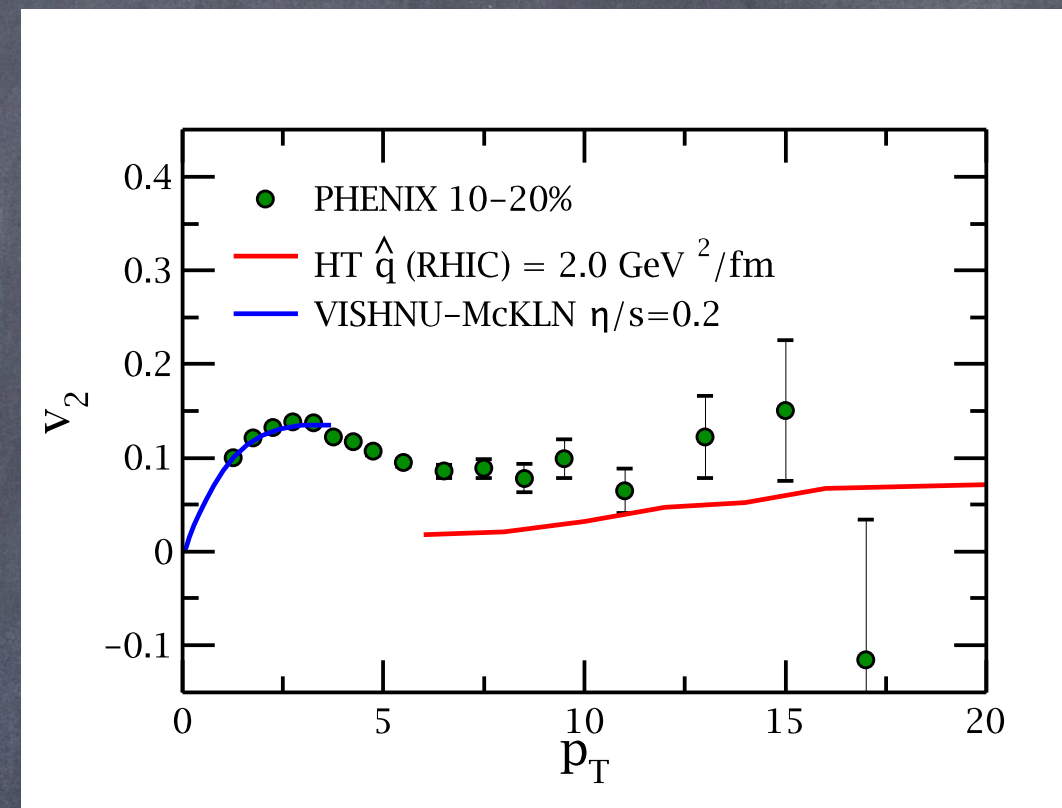
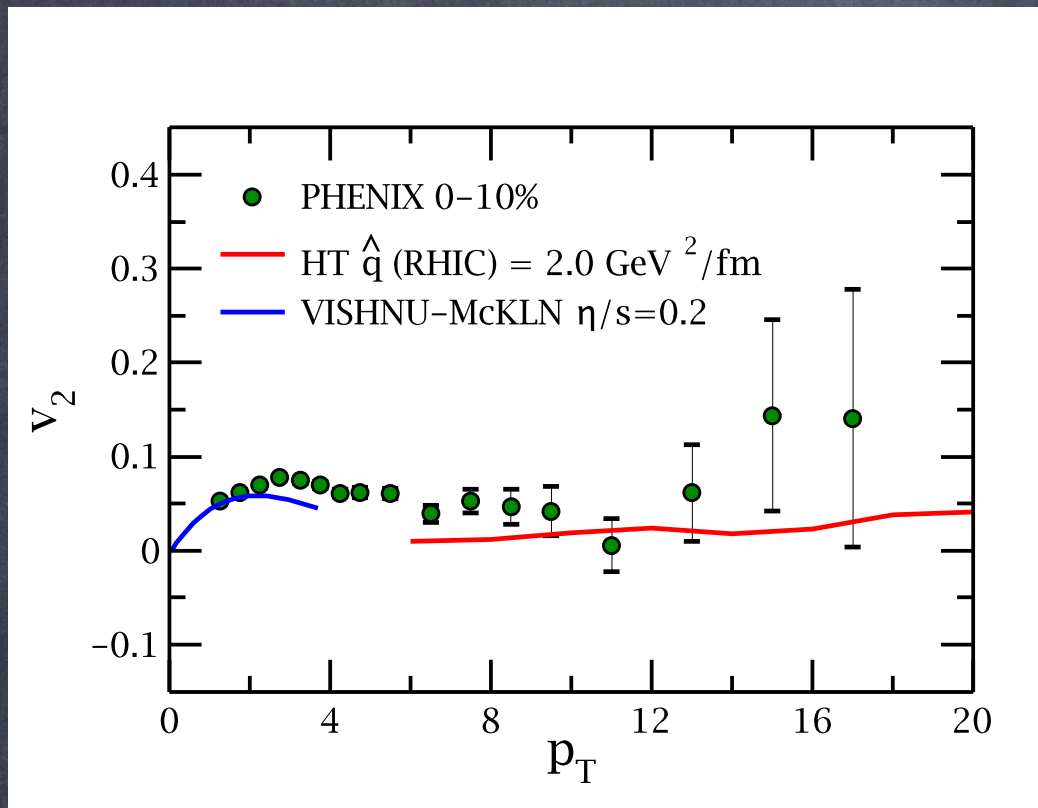
Tested in a 2+1D viscous hydro simulation



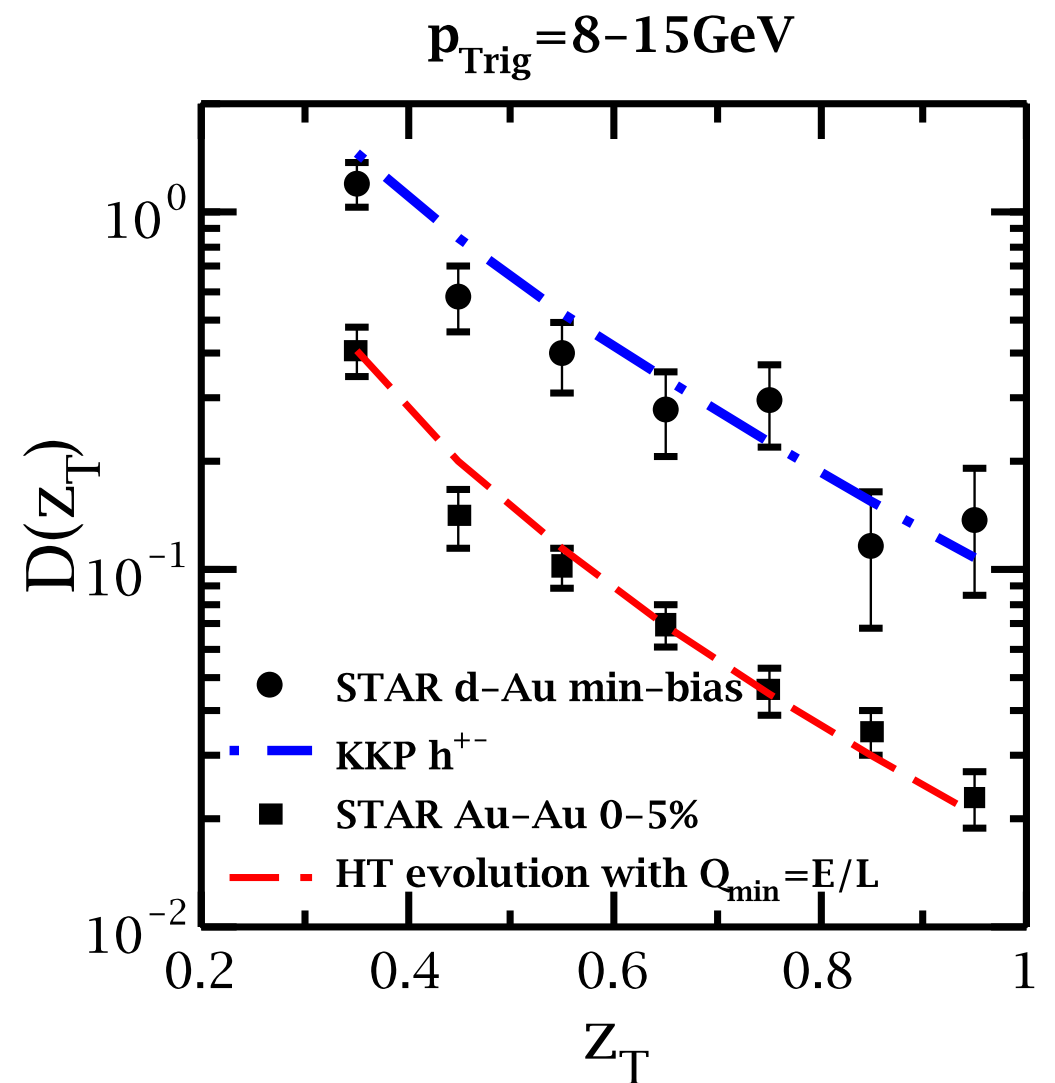
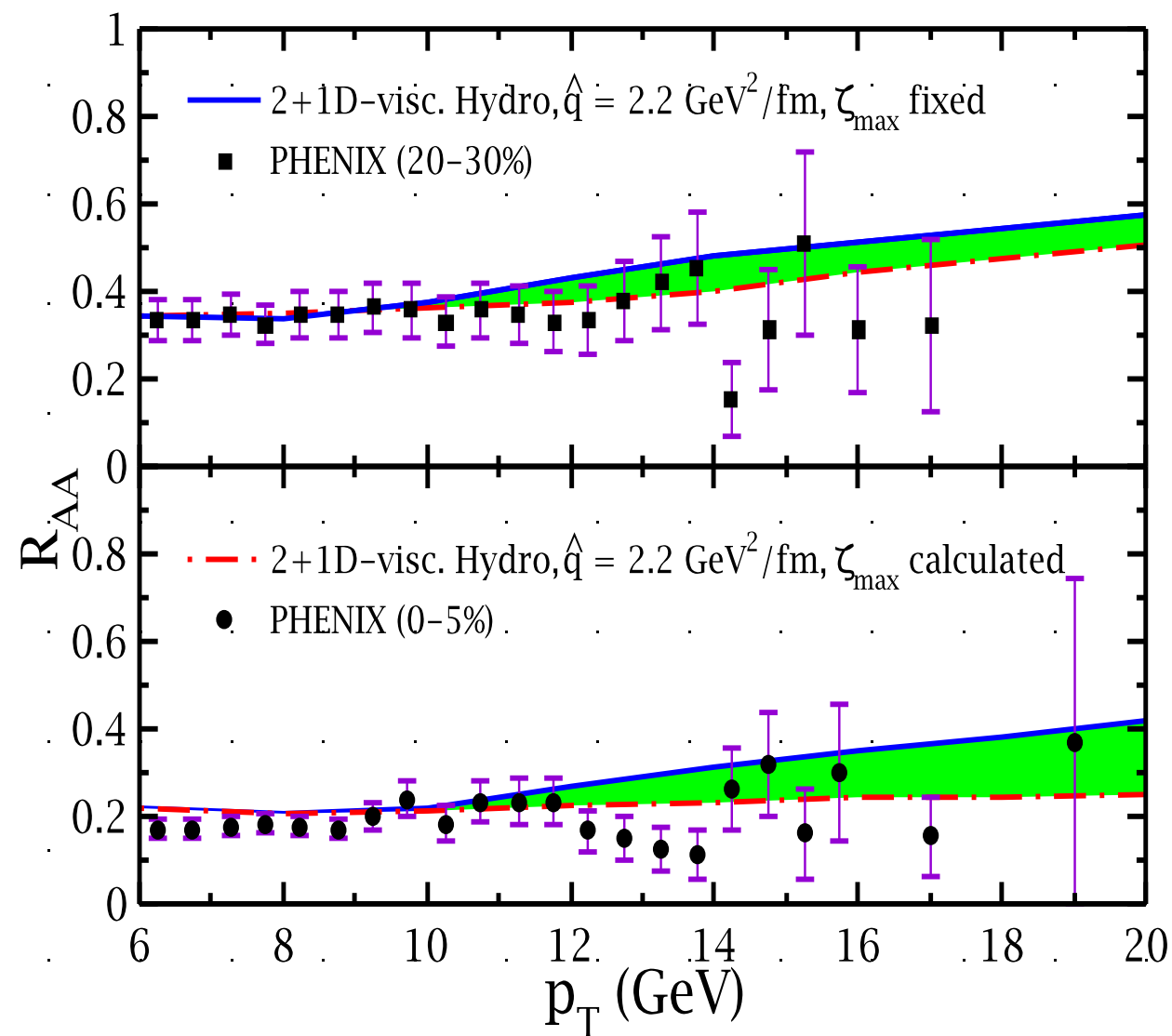
v_2 at LHC without a bump in \hat{q}/T^3



v_2 at RHIC without a bump in \hat{q}/T^3



Tested in parameter free extensions to away side



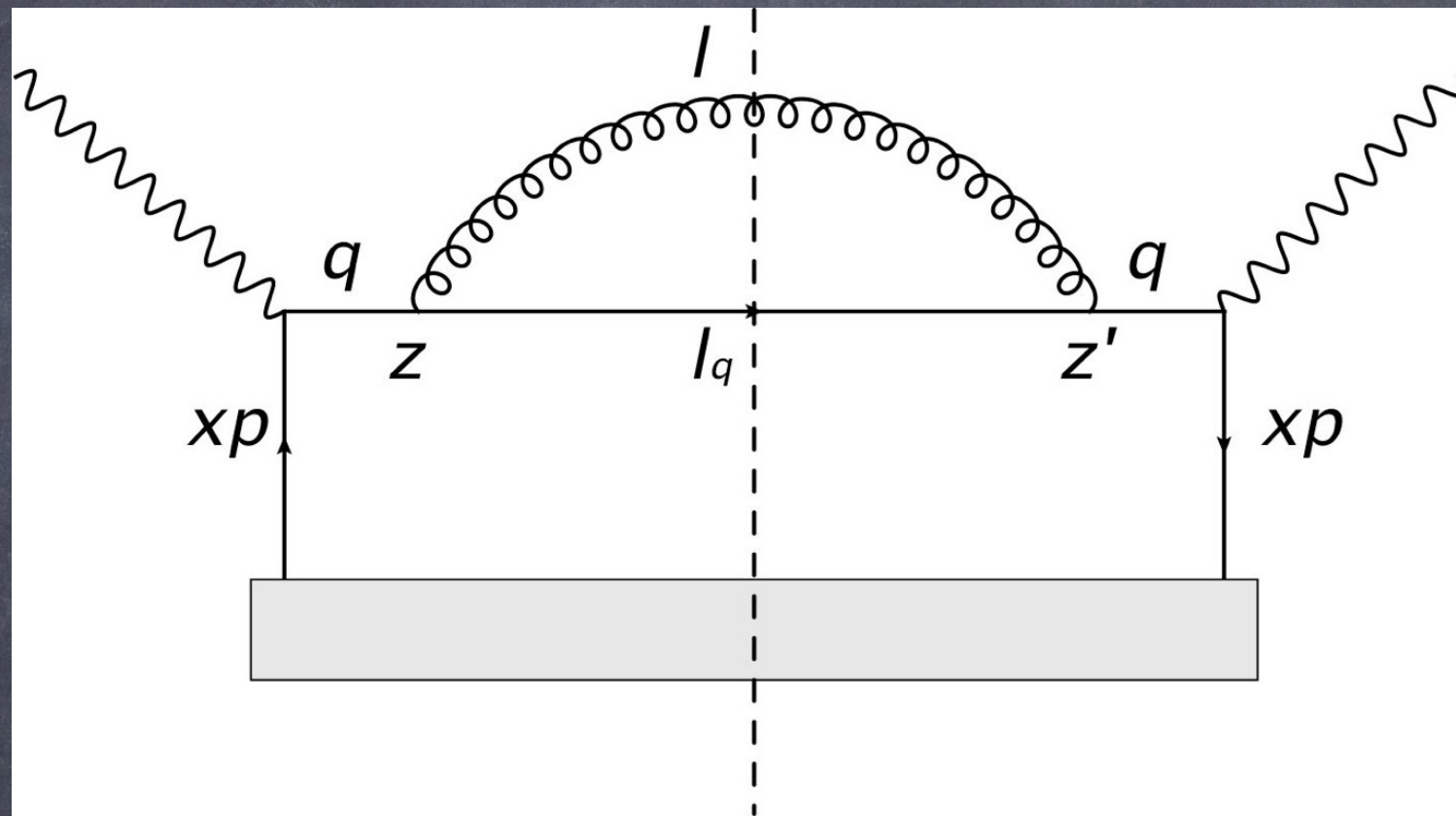
And extrapolated to LHC energies!

MATTER review

Main problem: Introducing distance into a DGLAP shower

No space-time in the usual Monte-Carlo showers

$$\bar{z} = \frac{z + z'}{2}$$



$$\delta z = z - z'$$

what is the role of z and z' ?

$$\int_0^\infty d^4 \bar{z} \exp [i(\delta q) \bar{z}] \quad \int d^4 \delta z \exp [i \delta z (l + l_q - q)]$$

δq is the uncertainty in q ,

How much uncertainty can there be ?

To be sensible: $\delta q \ll q$

we assume a Gaussian distribution around q^+

And try different functional forms of the width

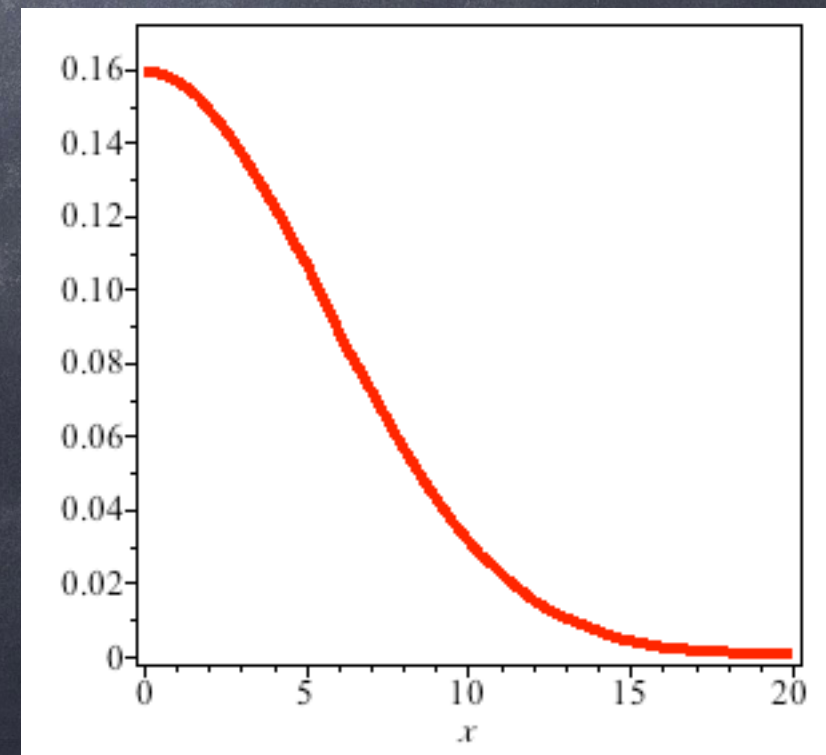
We set the form by insisting $\langle \tau \rangle = 2q^+/(Q^2)$

to obtain the z^- distribution only need to assume a δq^+ distribution

$$\rho(\delta q^+) = \frac{e^{-\frac{(\delta q^+)^2}{2[2(q^+)^2/\pi]}}}{\sqrt{2\pi[2(q^+)^2/\pi]}}$$

A normalized Gaussian with
a variance $2q^+/\pi$

FT gives
the following
distribution in
distance

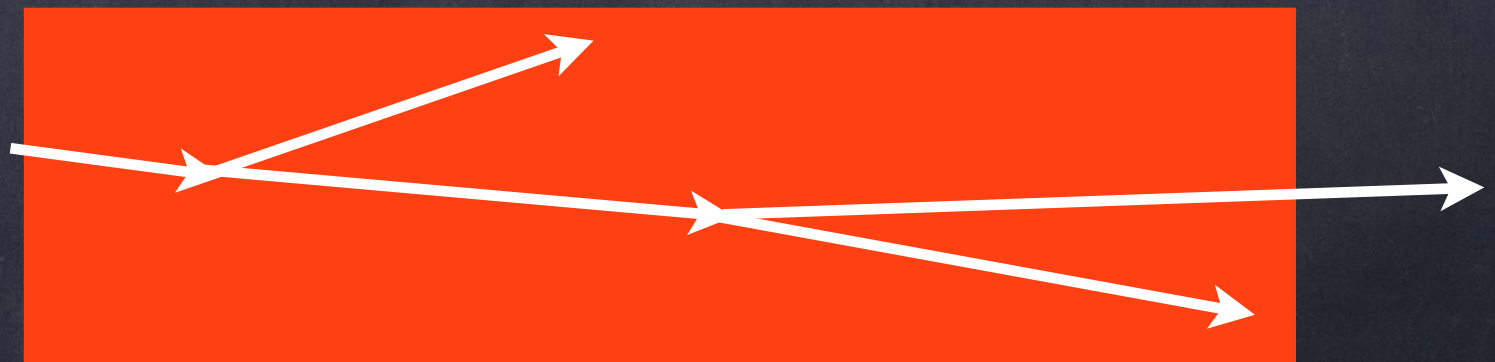


Constructing the Sudakov form factor

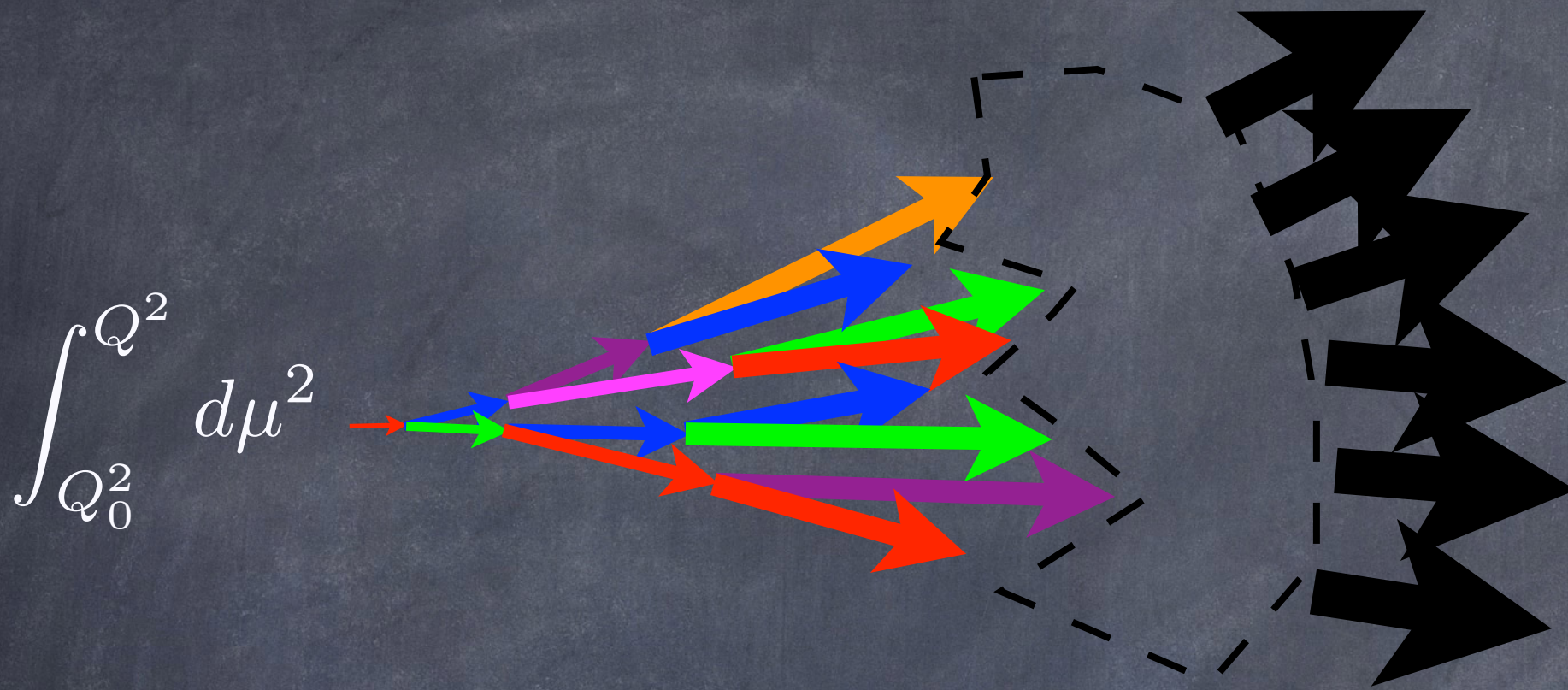
$$S_{\zeta_i^-}(Q_0^2, Q^2) = \exp \left[- \int_{2Q_0^2}^{Q^2} \frac{d\mu^2}{\mu^2} \frac{\alpha_S(\mu^2)}{2\pi} \right. \\ \left. \times \int_{Q_0/Q}^{1-Q_0/Q} dy P_{qg}(y) \left\{ 1 + \int_{\zeta_i^-}^{\zeta_i^- + \tau^-} d\zeta K_{p^-, \mu^2}(y, \zeta) \right\} \right]$$

$$K_{p^-, \mu^2}(y, \zeta) = \frac{2\hat{q}}{\mu^2} \left[2 - 2 \cos \left\{ \frac{\mu^2(\zeta - \zeta_i)}{2p^- y(1-y)} \right\} \right]$$

Valid as long as, $\frac{\hat{q}\tau}{\mu^2} \lesssim 1$



Mass, Virtuality and Scale



Q^2 is the scale,

μ is the mass of jet = virtuality of first parton

virtuality drops as more shower partons produced

Reconstructed mass: $M = \sqrt{p^2 - p_T^2 - p_z^2}$

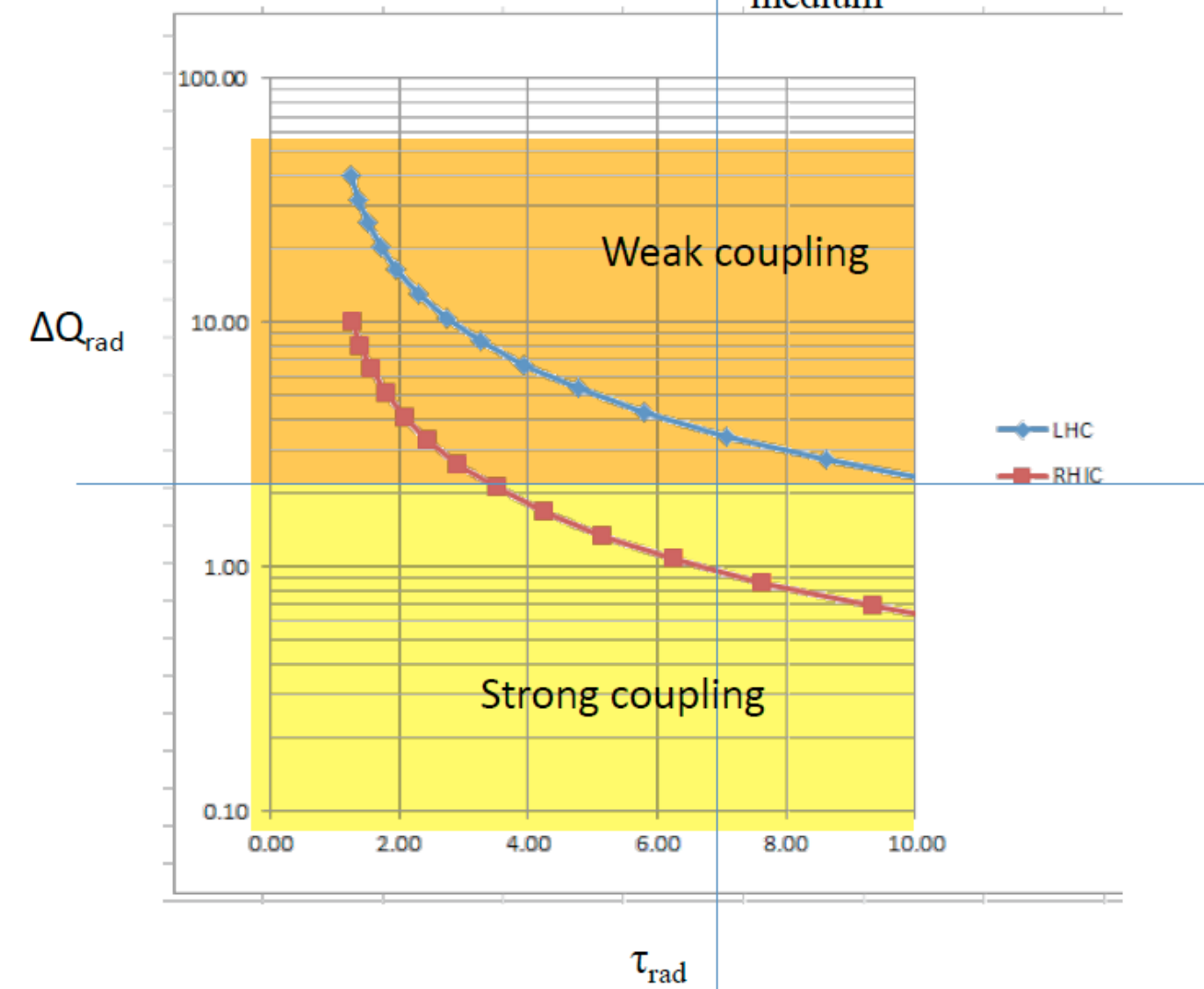
$$\vec{p}_T = \sum_{i=1}^n \vec{p}_{T_i}, \quad p_T = |\vec{p}_T|, \quad p_{T_i} = |\vec{p}_{T_i}| \quad p_z = \sum_{i=1}^n p_{T_i} \sinh \eta_i, \quad p = \sum_{i=1}^n p_{T_i} \cosh \eta_i.$$

What happens to the leading parton in vacuum/medium?

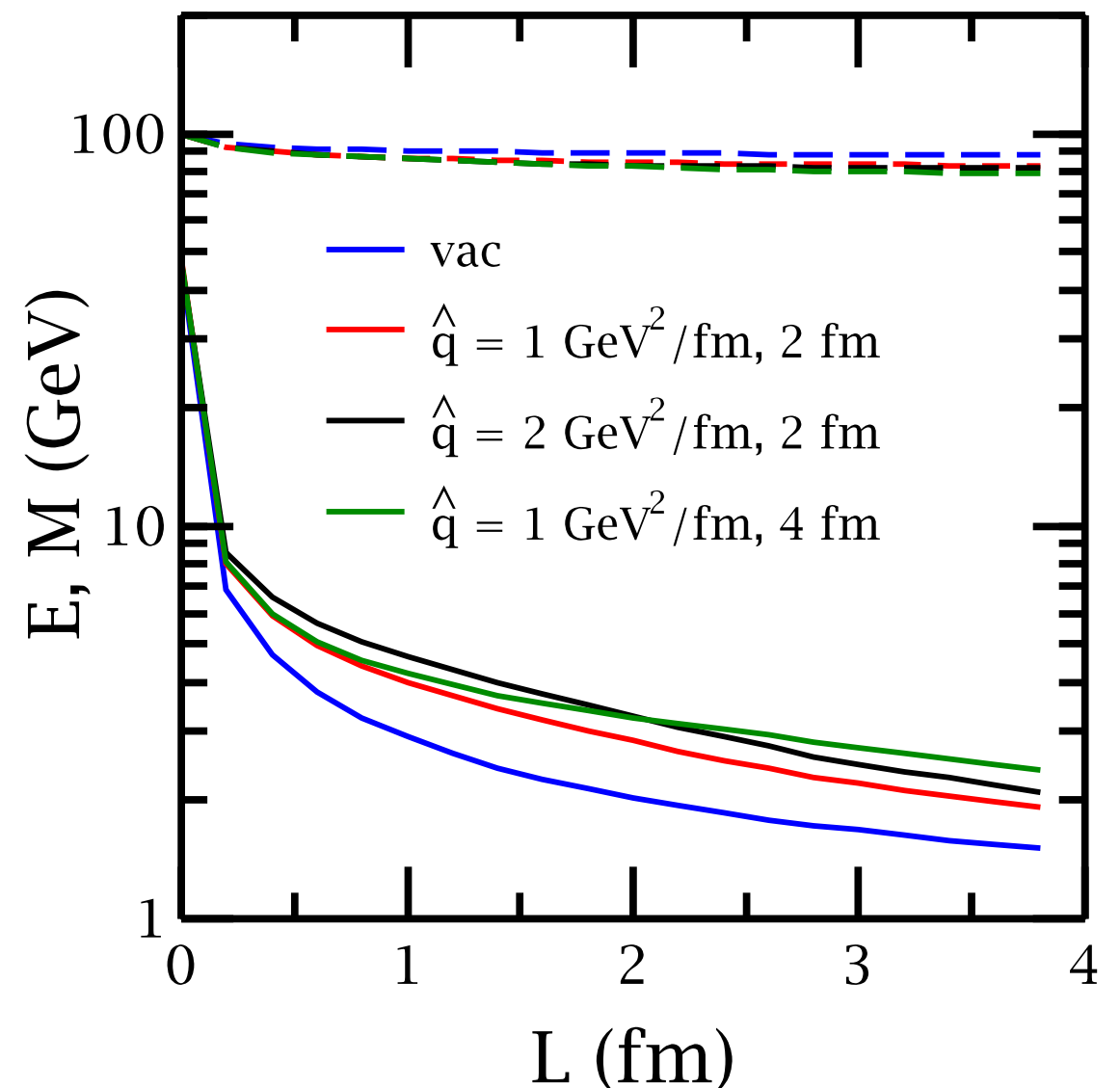
Virtuality or mass drops much more quickly than Energy

From talks by
B. Mueller and R. Seto

→
 τ_{rad} out of
medium



static Medium

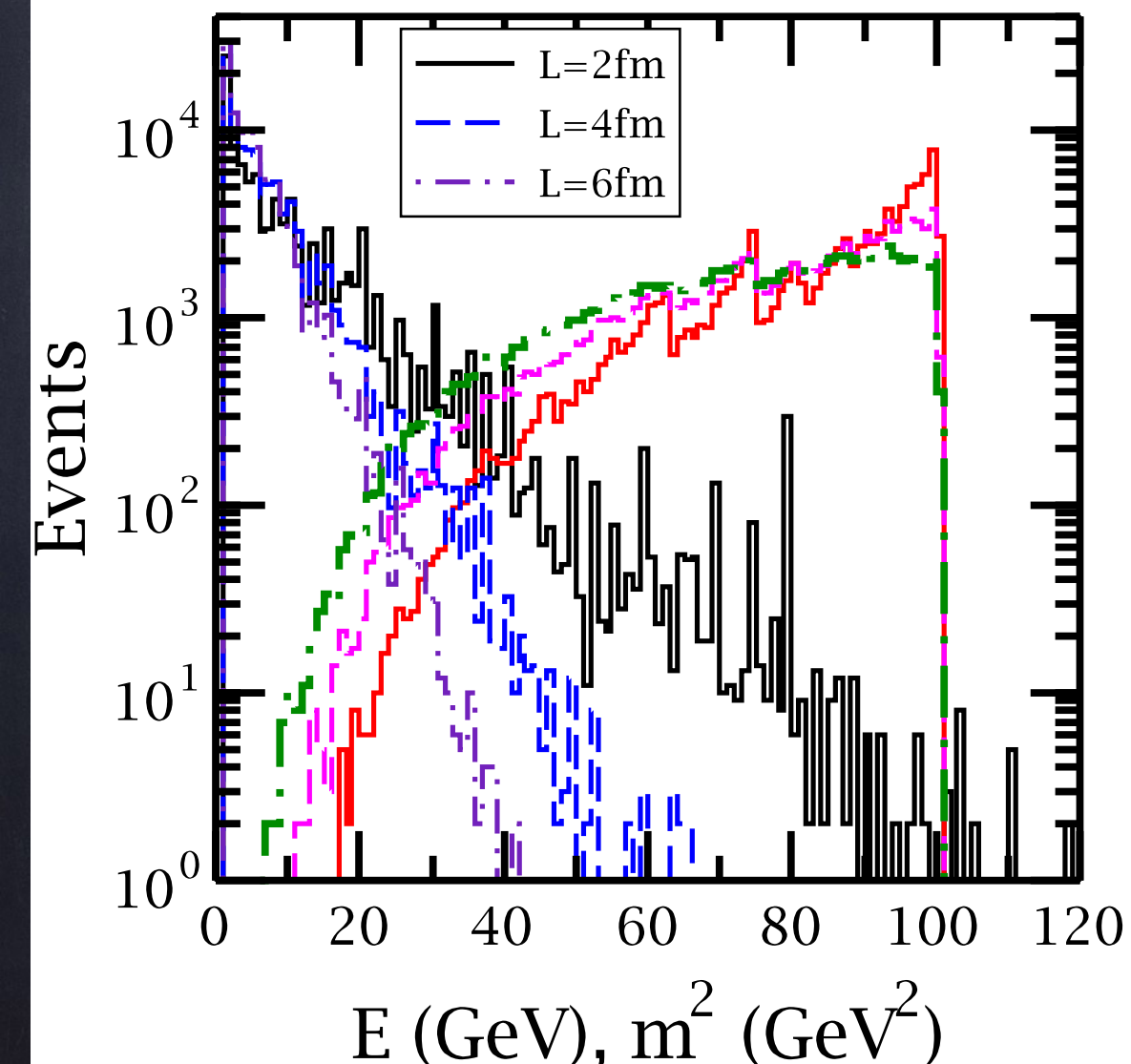


Note: once partons drop to 1GeV, no further splitting occurs

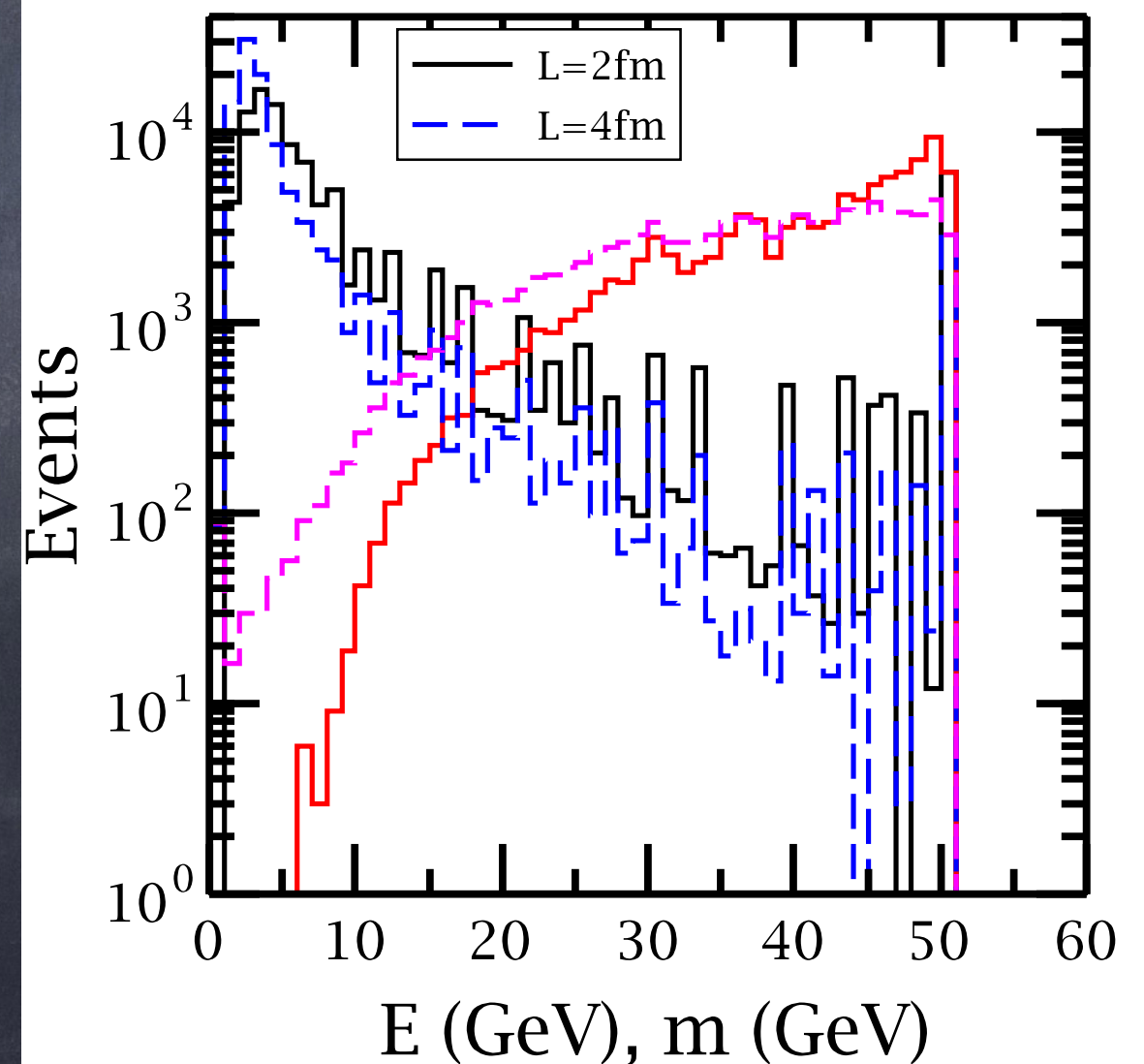
What happens to the leading parton in vacuum/medium?

How about the distributions?

E=100GeV, 100K events



E=50GeV, 100K events



Distribution of virtuality from initial hard scattering

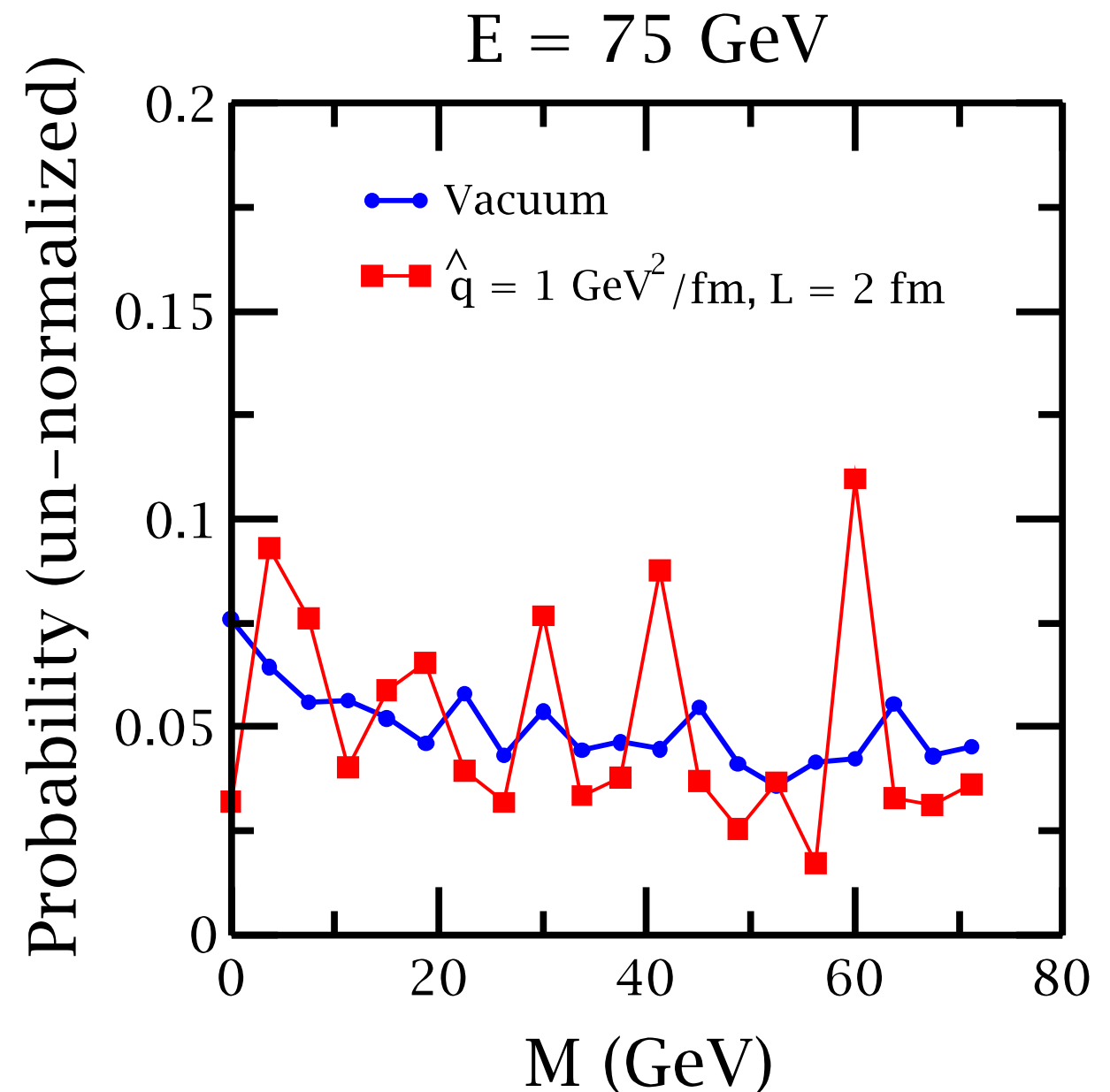
Jagged distribution due to fewer events

$M < 1 \text{ GeV}$ clubbed with 0 GeV point

For $E = 75 \text{ GeV}$

scale = max. virtuality
= 75 GeV

mass distribution of leading parton From
MATTER with no initial state (100K events)



Distribution of virtuality from initial hard scattering

Jagged distribution due to fewer events

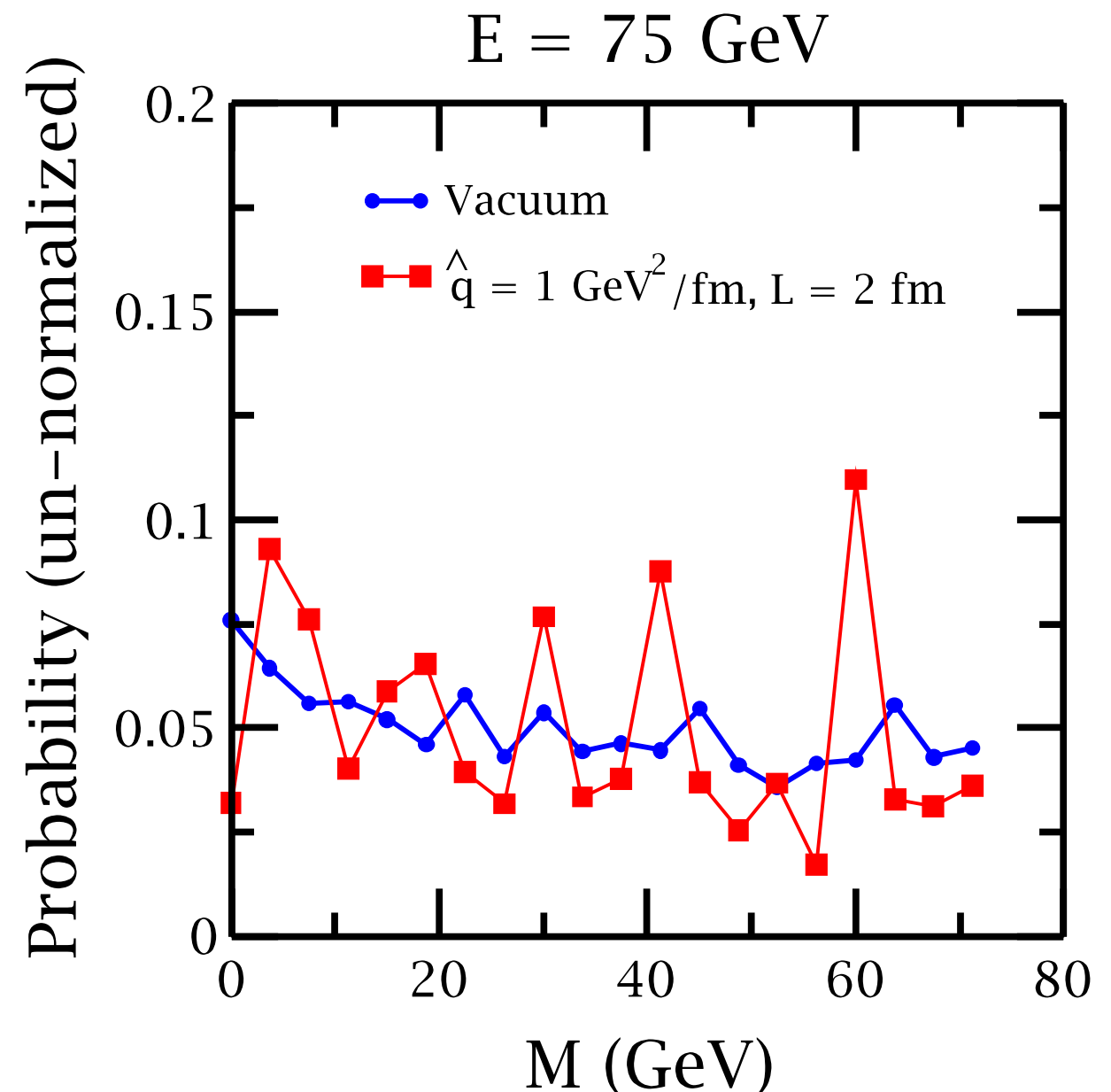
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vs. Virt distribution of recon. Jet with
 $q = 1 \text{ GeV}^2/\text{fm}$ and $L = 2 \text{ fm}$

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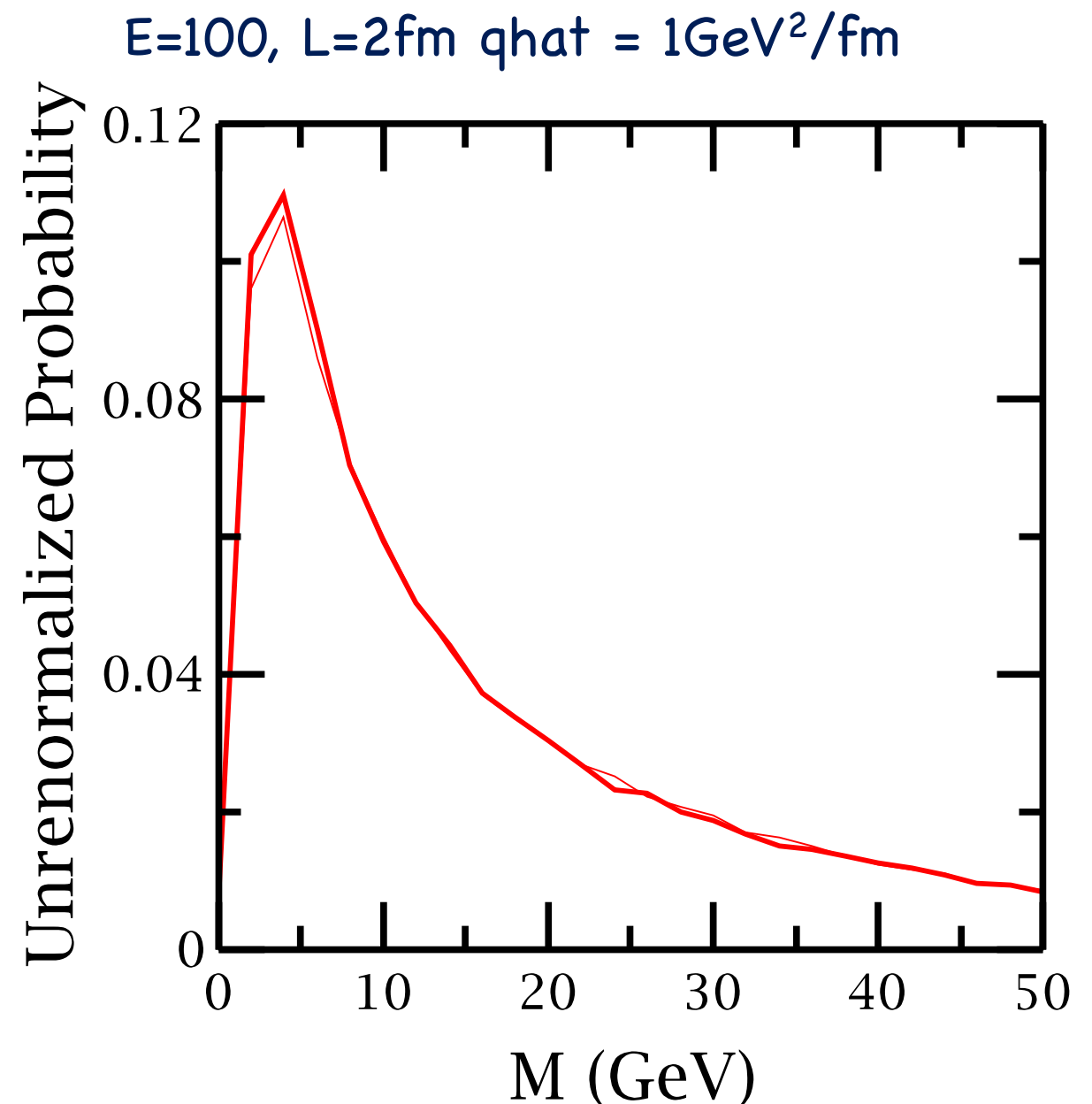
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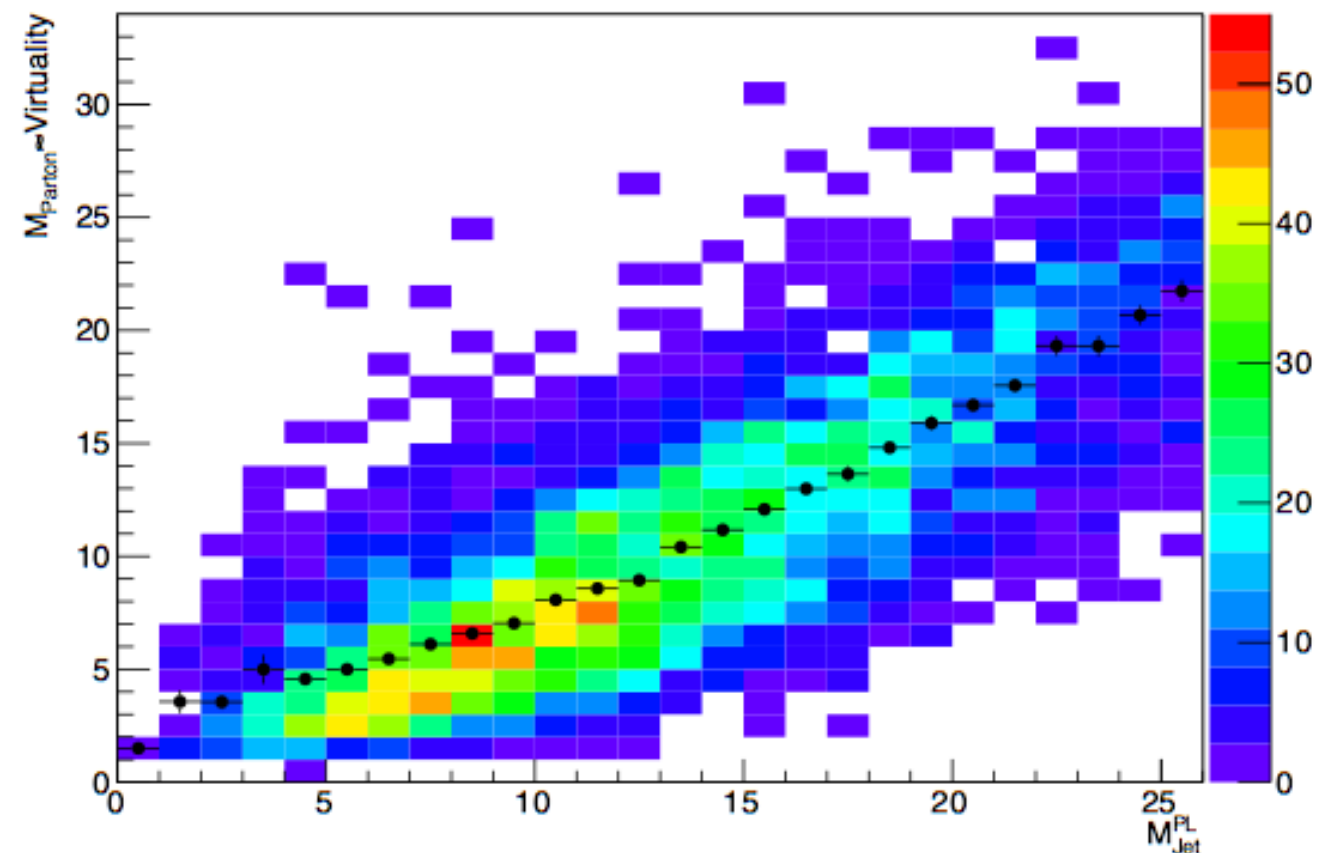
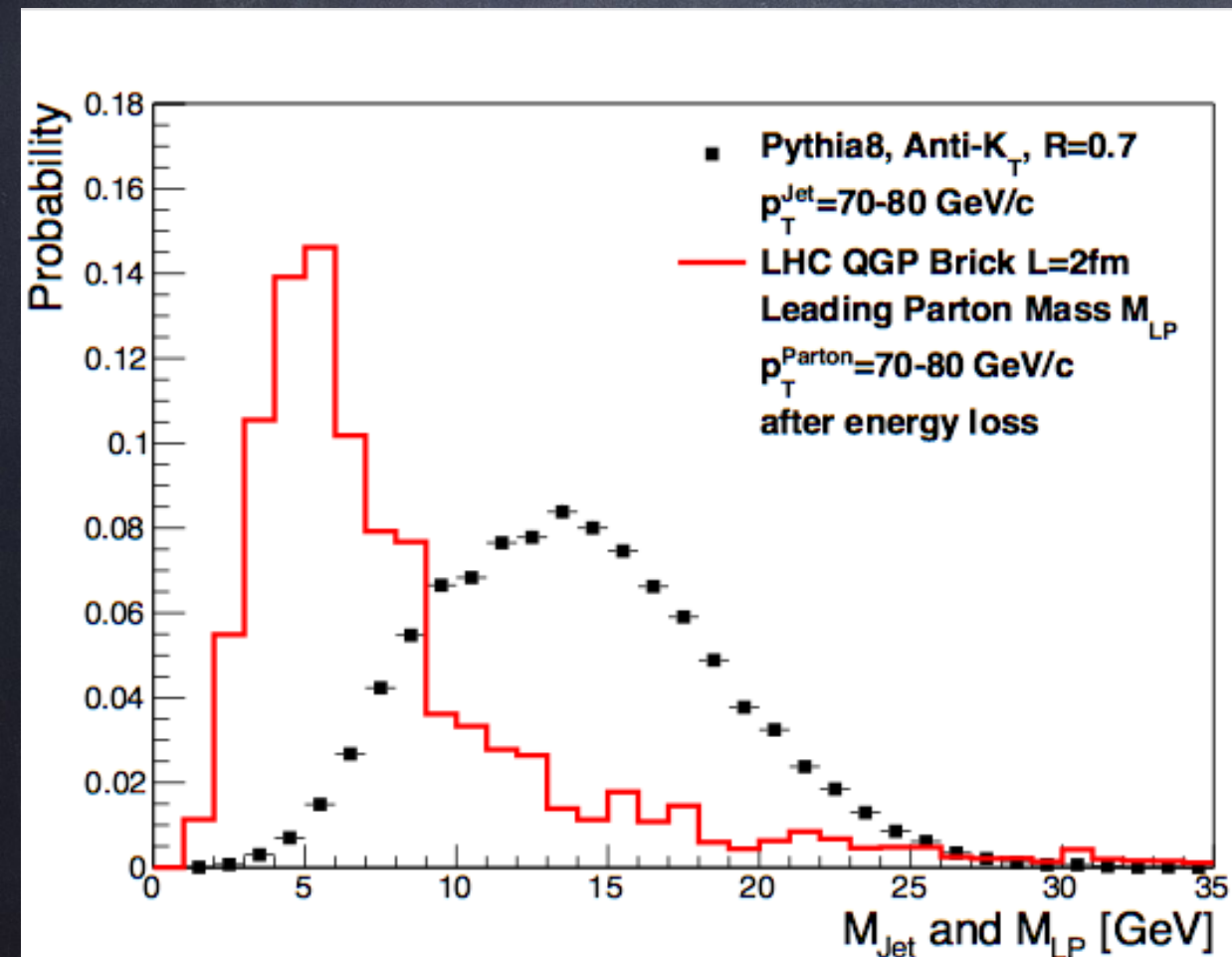


The large difference is also present for PYTHIA jets

Due to presence of hard scattering and initial state the mass distribution in PYTHIA has a peak
Here we use reconstructed mass,
which compares well with actual mass

Mass distribution of reconstructed jet in
PYTHIA vs in-medium reconstructed jet

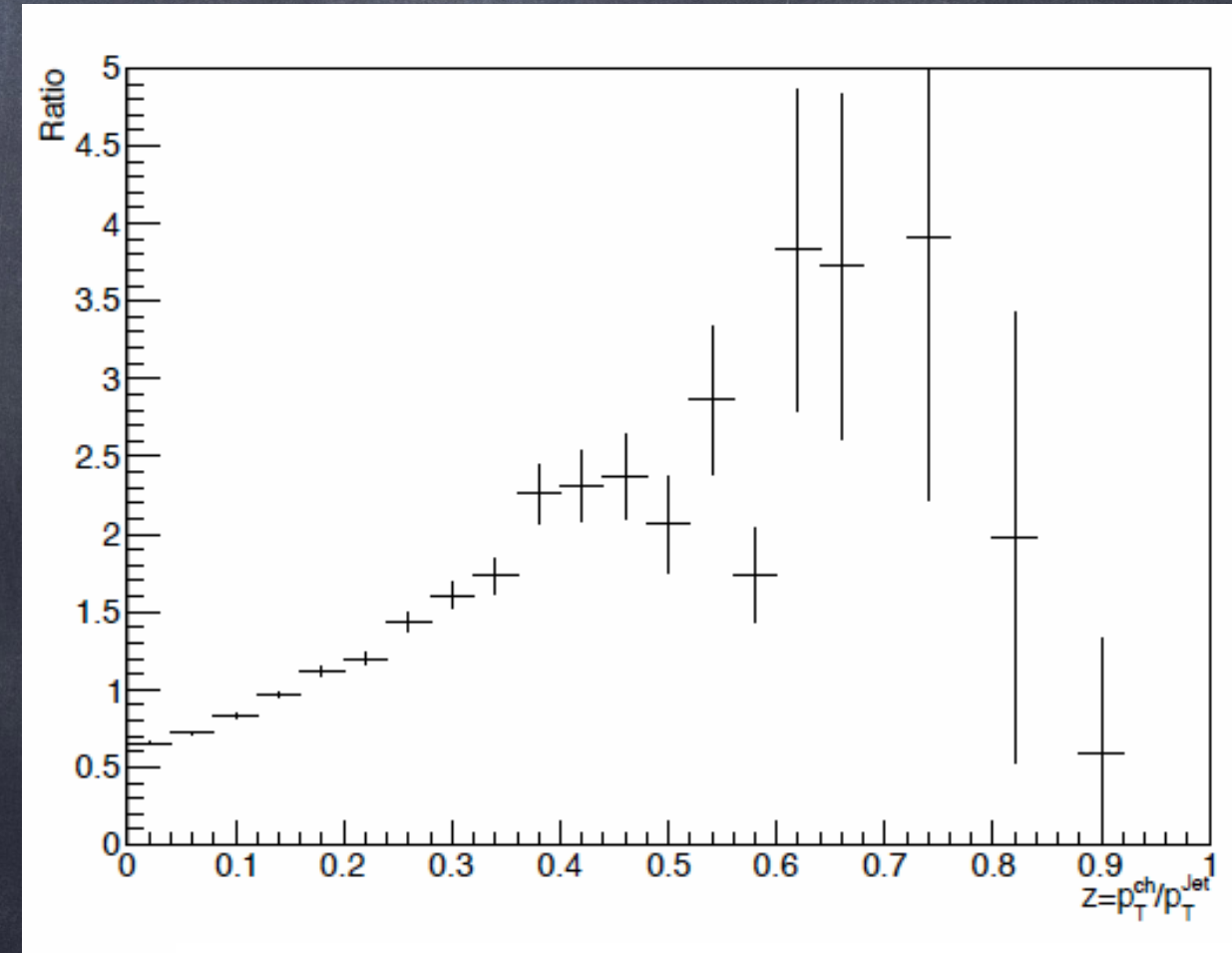
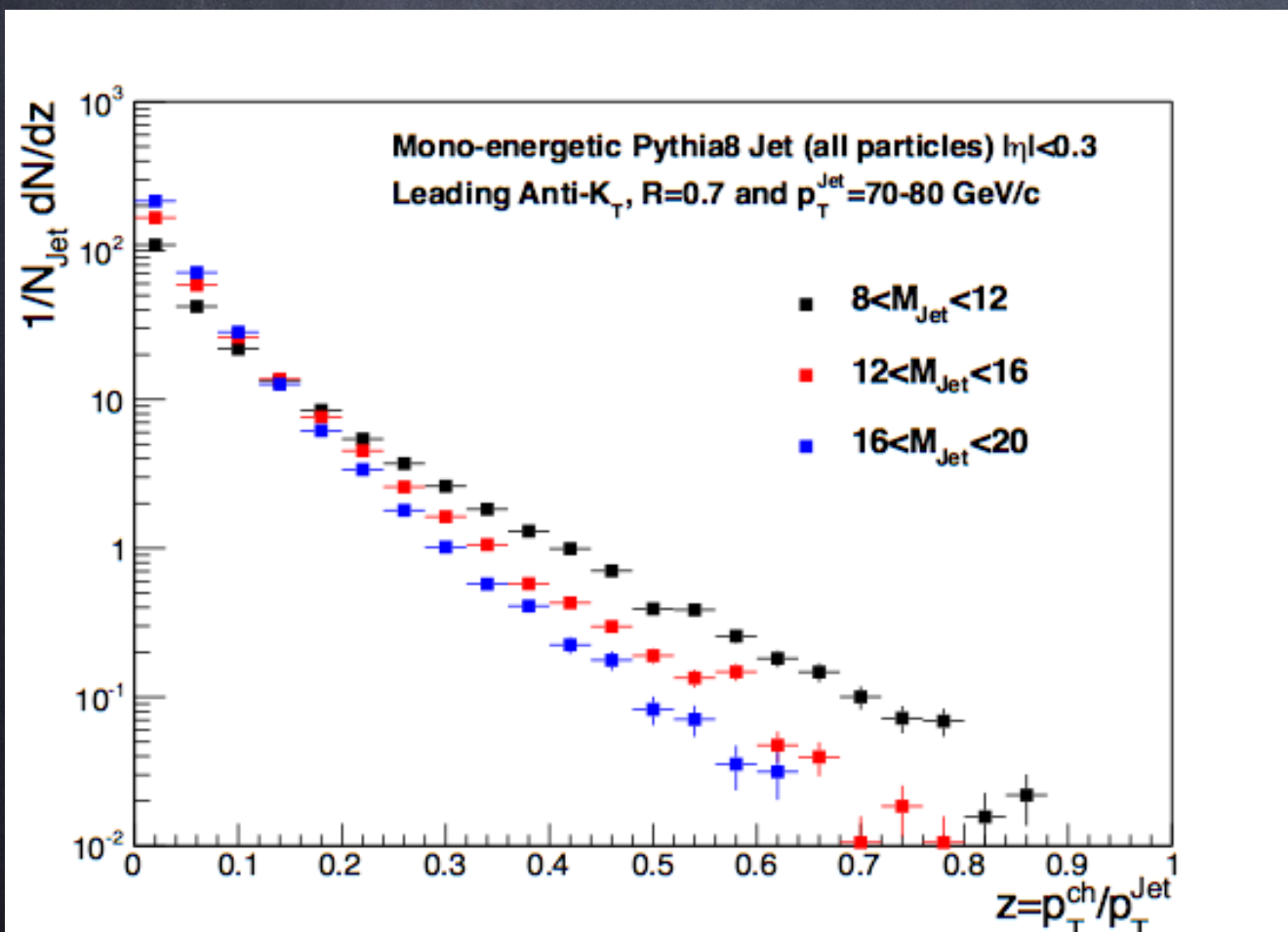
Reconstructed mass is very closely
related to actual mass in PYTHIA



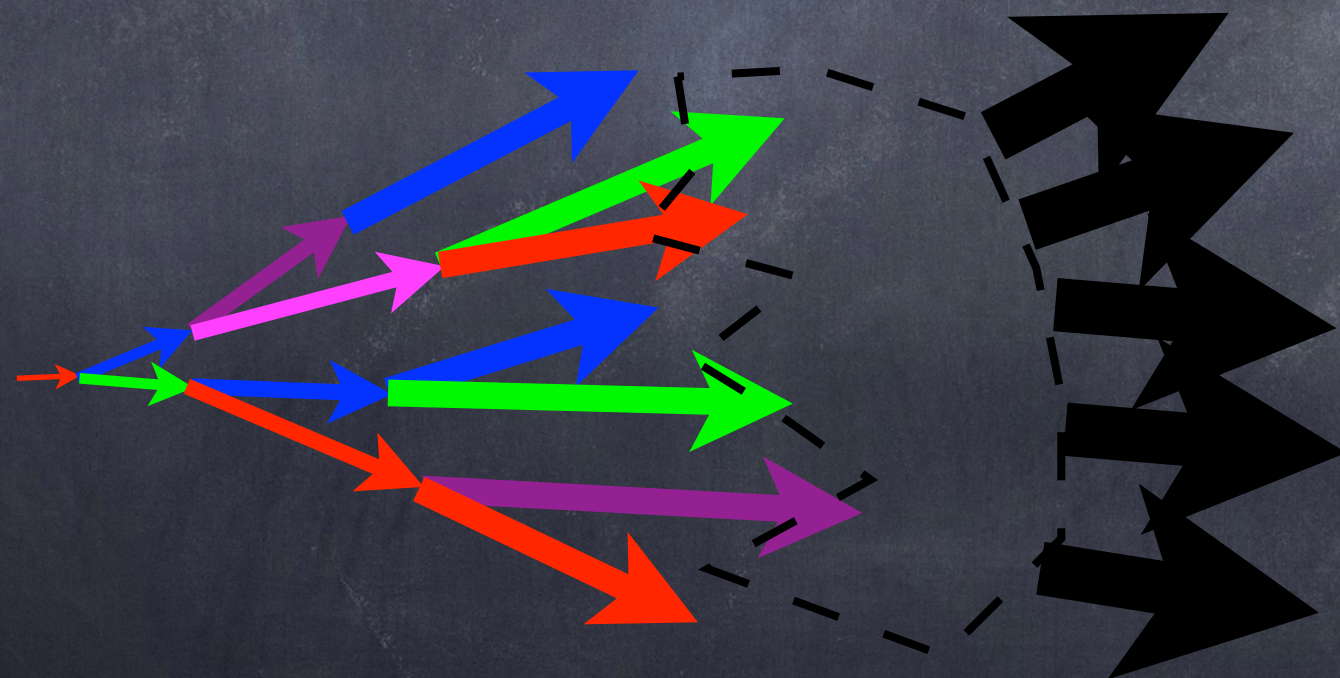
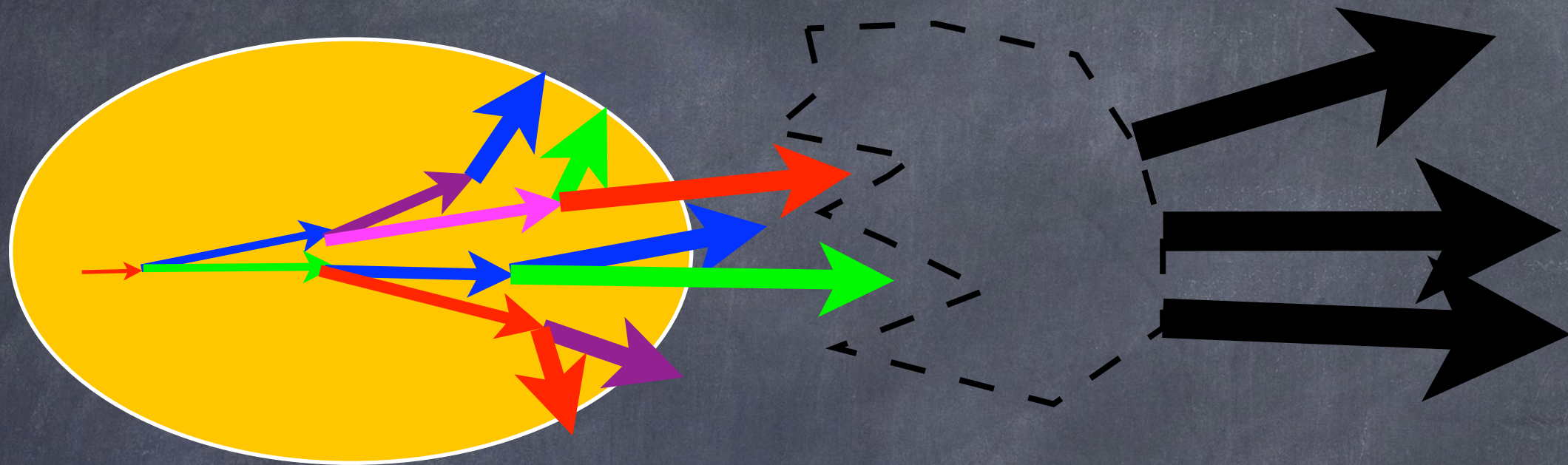
Reconstructed mass as a place holder for jet mass

The jet fragmentation function for a given mass range is very sensitive to that mass range

sensitivity very similar to the scale dependence of fragmentation function



We are dividing two very different objects



Semi conclusions

While leading partons lose some energy in heavy-ion collisions

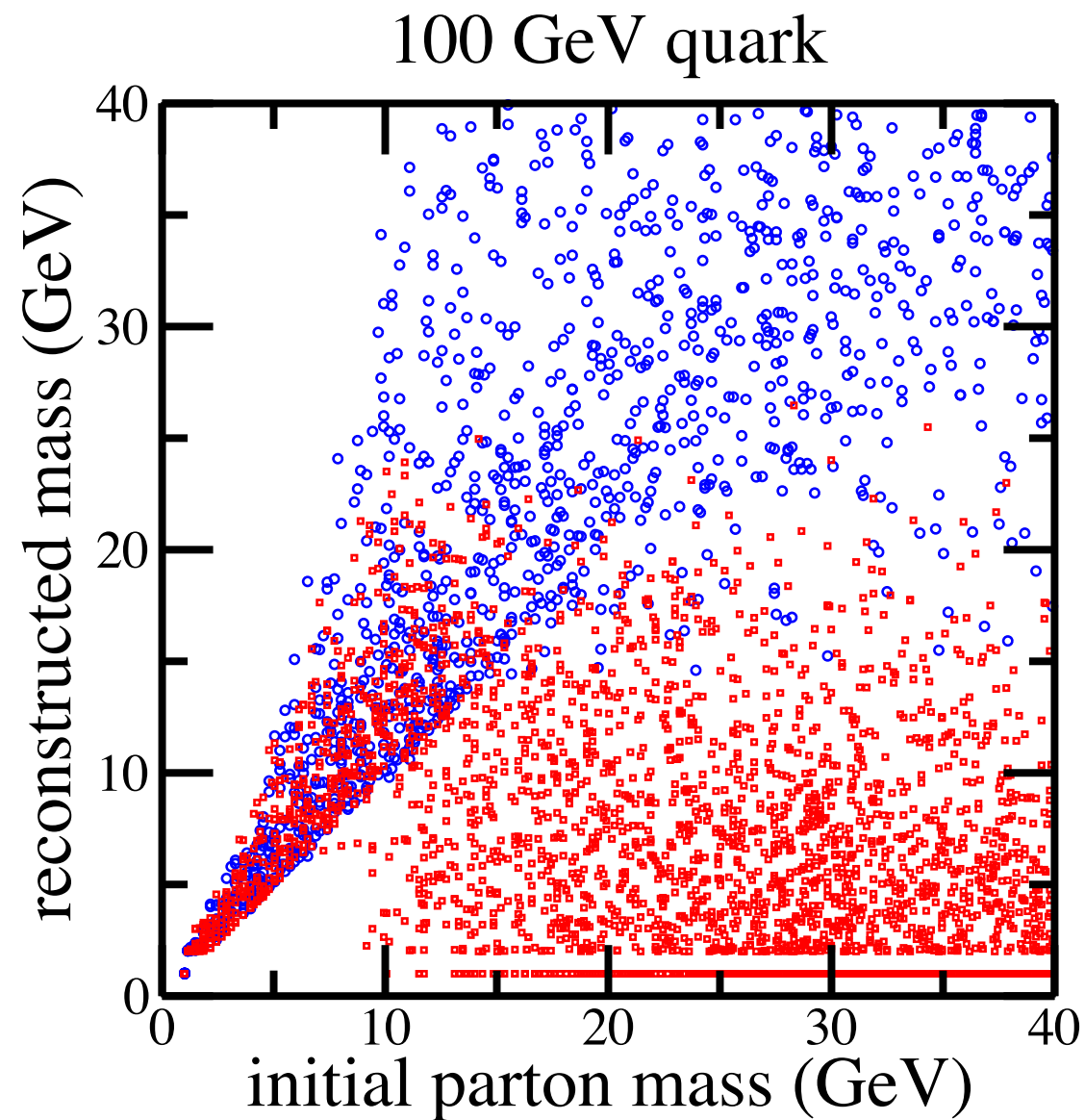
They lose a lot of mass in the course of radiation

If many of the soft partons are deflected in the medium, we are left with a mass-depleted jet

The mass of the depleted jet has a narrow distribution

This should be a noticeable effect if mass or some approximation to mass could be measured.

Reconstruction does have an effect on
extracted mass of jet.



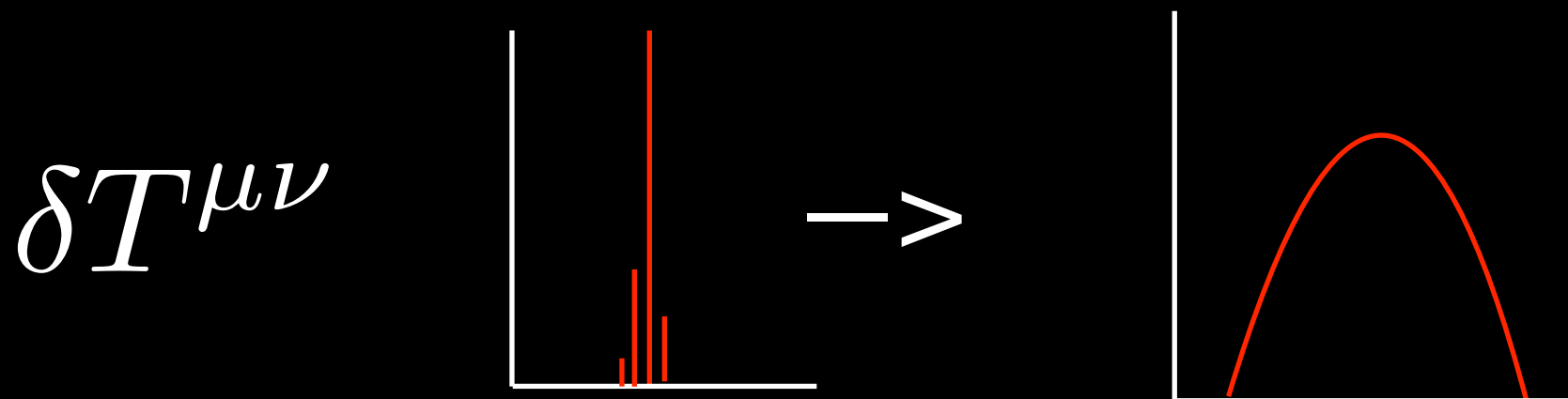
In general, 2 kinds of transport coefficients

Type 1: which quantify how the medium changes the jet

$$\hat{q}(E, Q^2) \qquad \hat{q}_4(E, Q^2) = \frac{\langle p_T^4 \rangle - \langle p_T^2 \rangle^2}{L} \dots$$

$$\hat{e}(E, Q^2) \qquad \hat{e}_2(E, Q^2) = \frac{\langle \delta E^2 \rangle}{L} \qquad \hat{e}_4(E, Q^2) = \frac{\langle \delta E^4 \rangle - \langle \delta E^2 \rangle^2}{L} \dots$$

Type 2: which quantify the space-time structure of the deposited energy momentum at the hydro scale



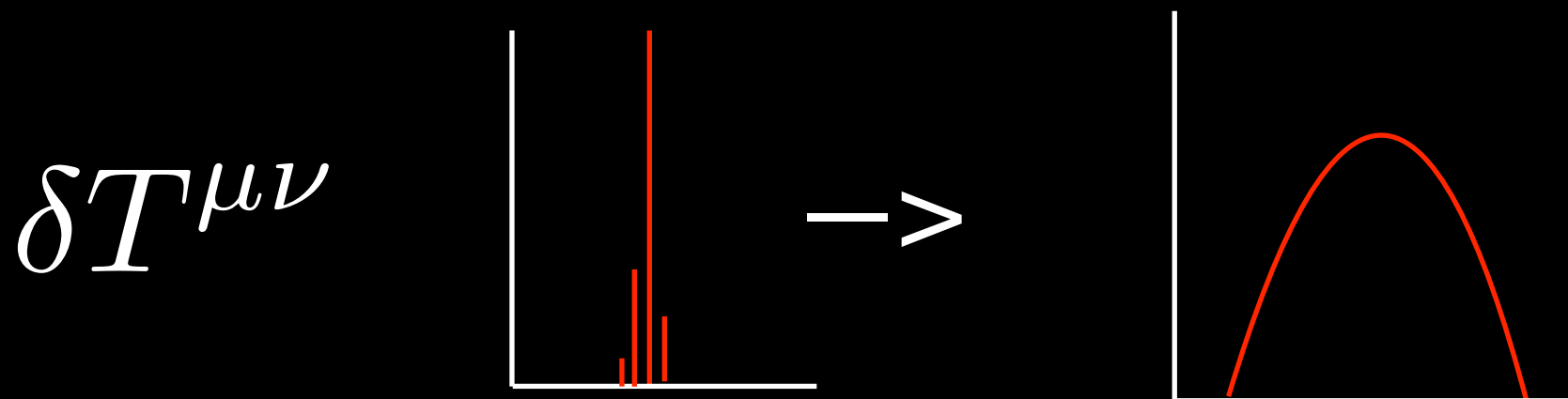
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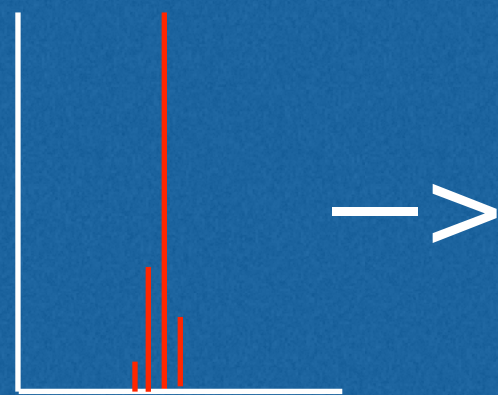
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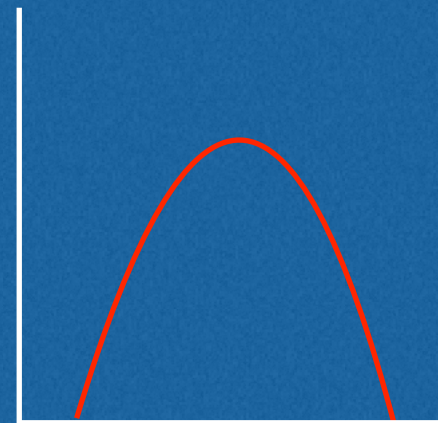
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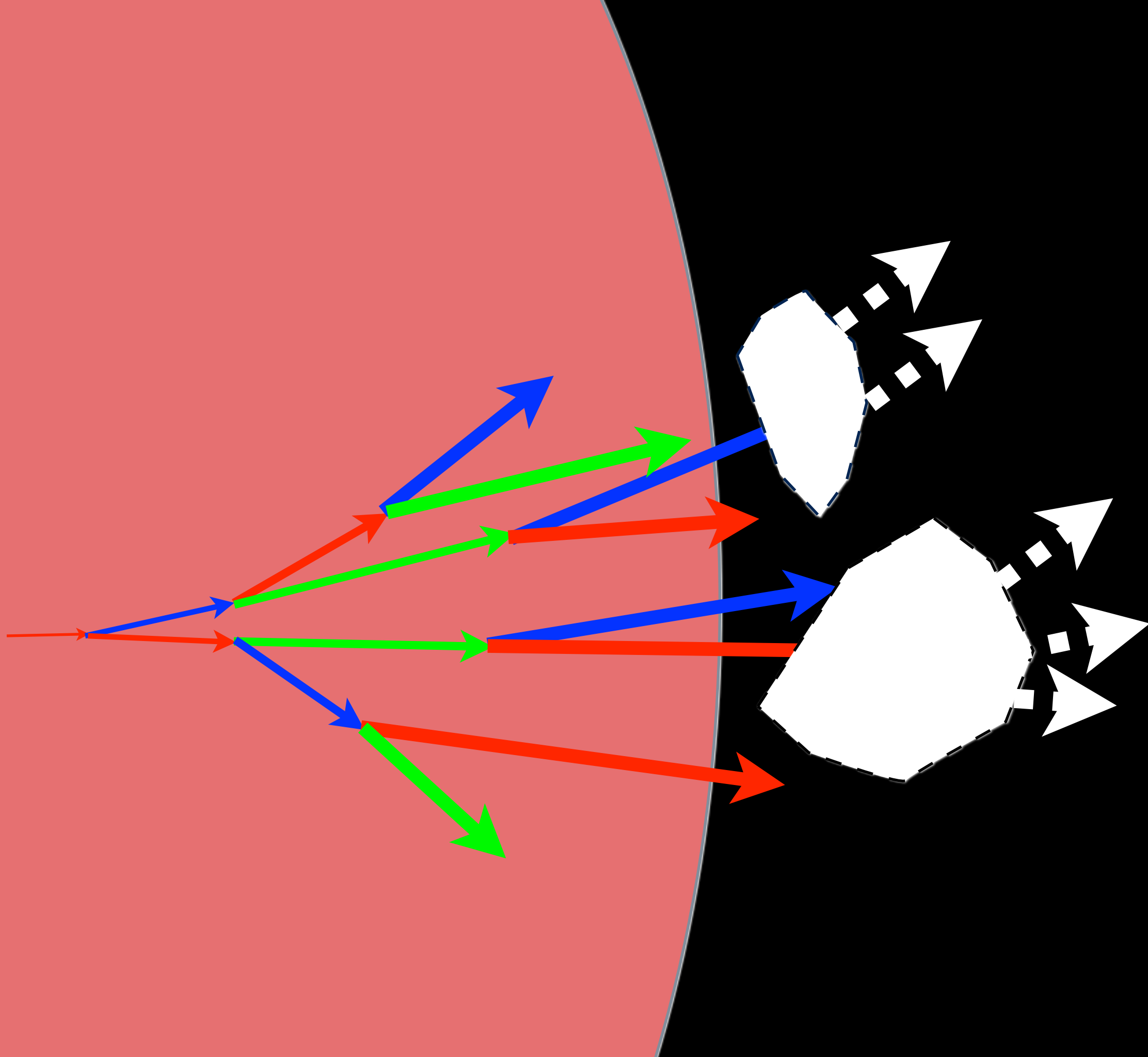
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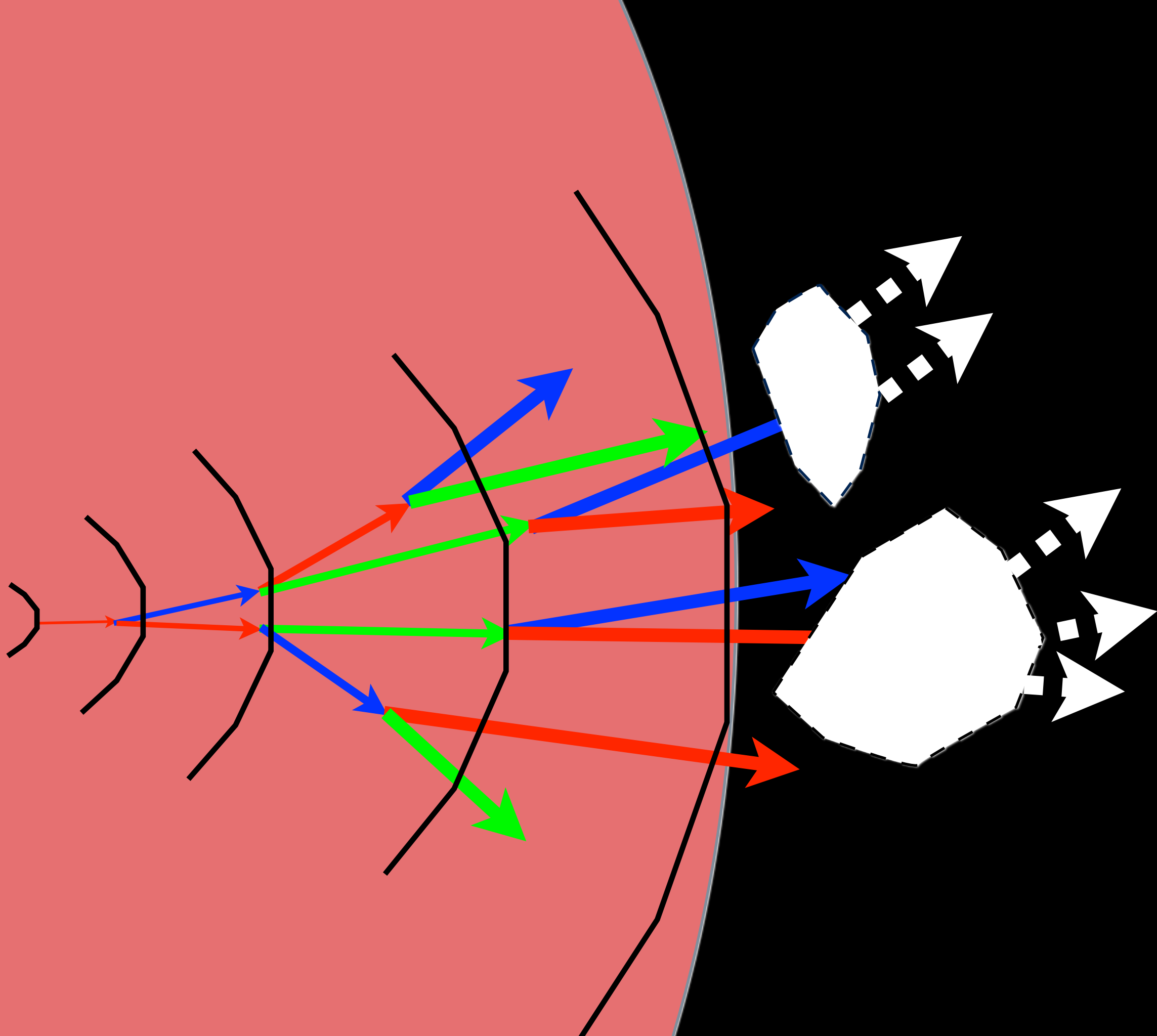
$\delta T^{\mu\nu}$



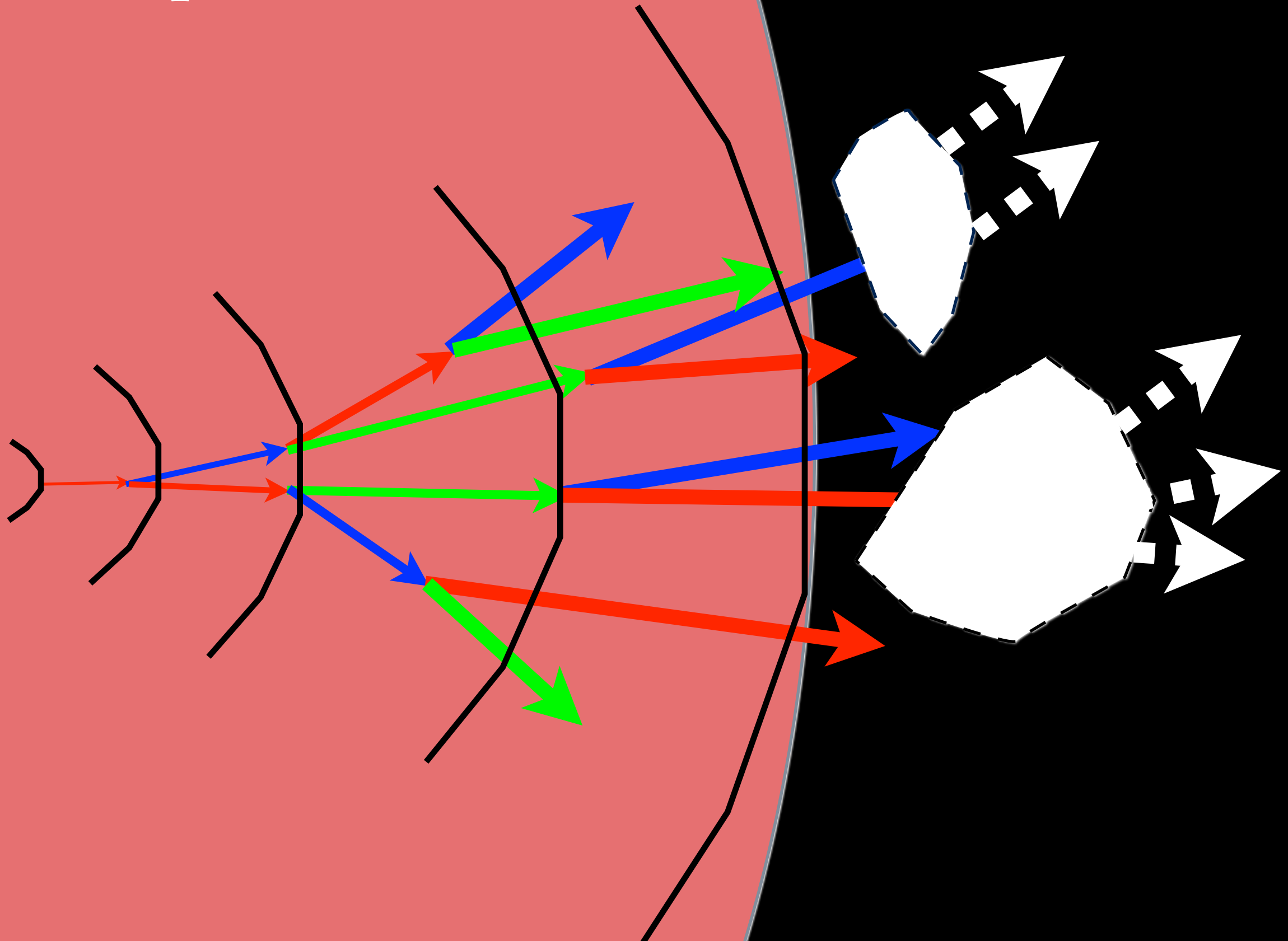
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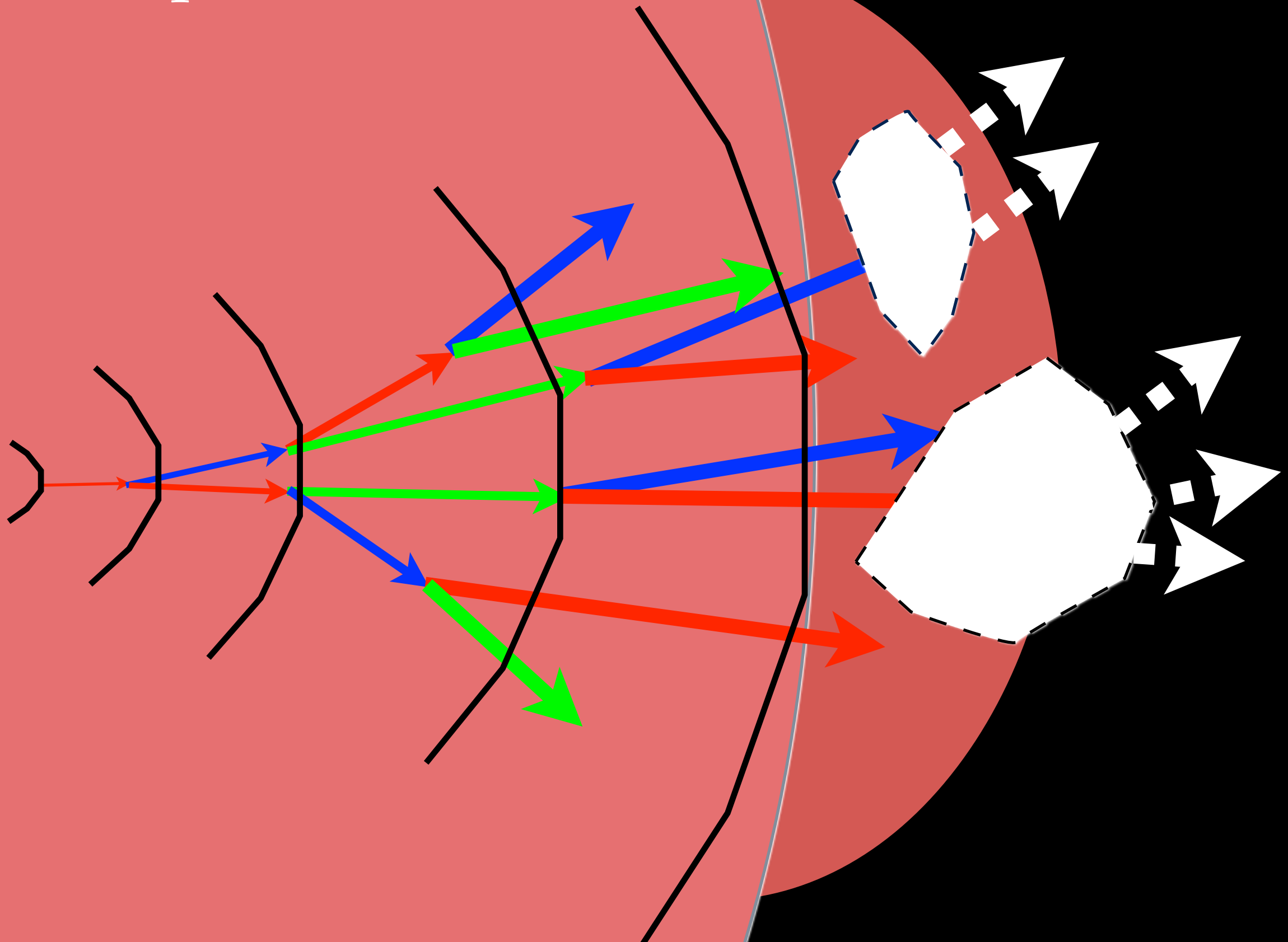




The modification of the underlying medium depends on these new coefficients



The modification of the underlying medium depends on these new coefficients



Conclusions II

Strong coupling input needed in event generators

Not just to complete the shower process

But to describe the redistribution of energy in jet cone that will effect the mass.